

Dynamical Impacts of convection and stochastic approaches

by

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ECMWF/Met Office

Outline

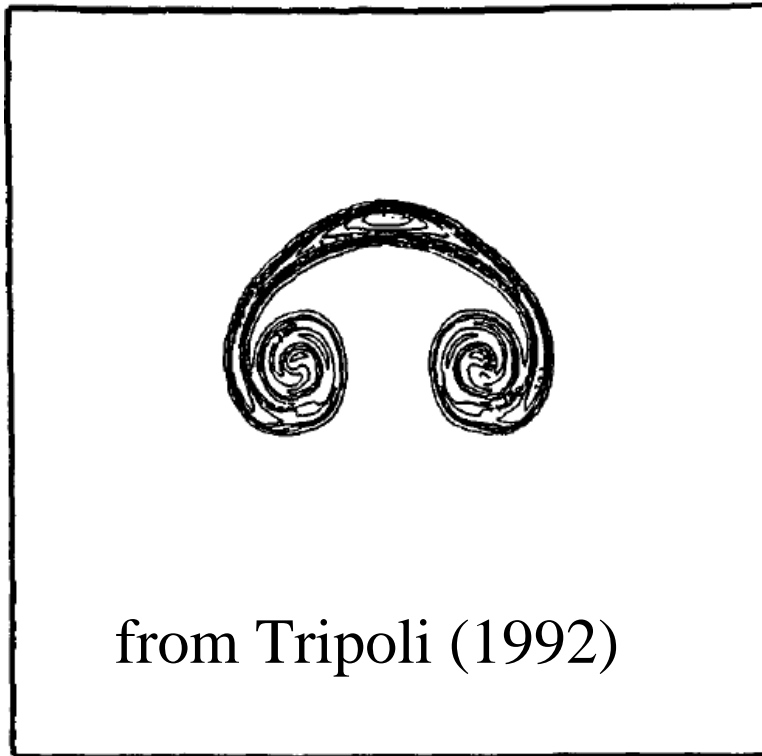
- **Dynamical processes in clouds and interactions with the rotating, stratified environment**
- **Generalized parcel model and balanced flow adjustment – PV generation and NWP impacts**
- **‘Big-domain’ tropical convection simulation – convectively-coupled tropical waves and statistical properties of convective forcing**
- **Stochastic parametrization and kinetic energy backscatter**

Convection conceptions



Mixing at cloud boundaries

- Horizontal gradients of buoyancy cause the baroclinic generation of vorticity to be concentrated in cloud boundaries making them unstable.
- Cloud droplet evaporation causes internal downdraughts



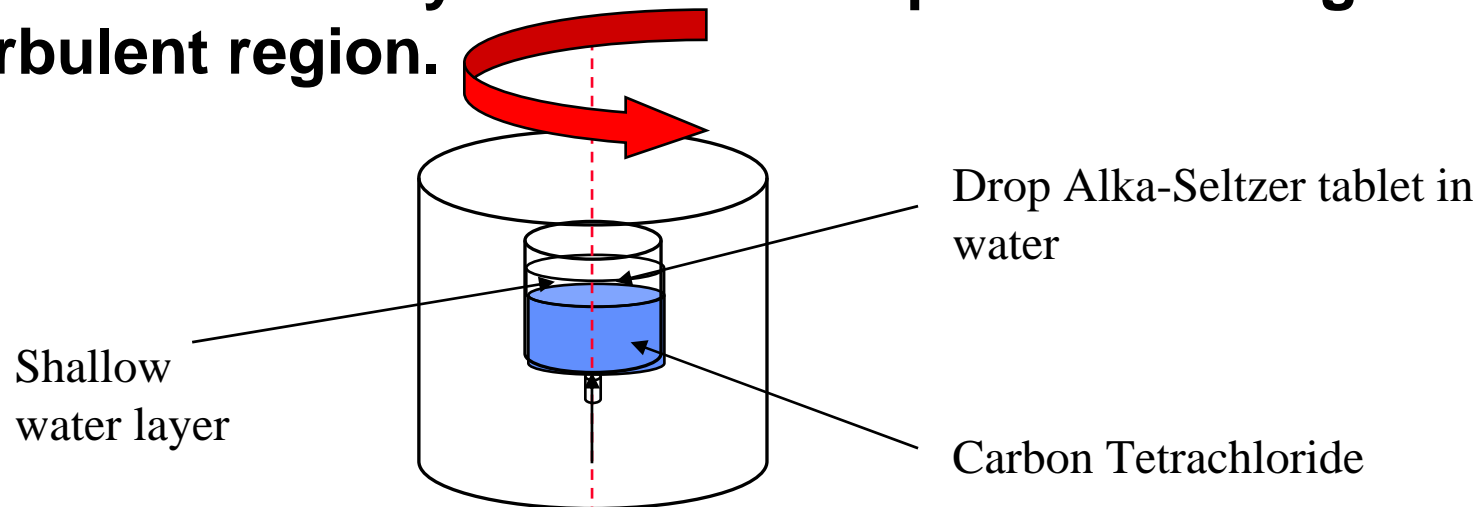
➔ **Turbulent clouds**

Vorticity production at thermal boundaries winds up into a double Swiss Roll

Interplay between background rotation and convection - the Alka-Seltzer experiment



- **Richard Scorer's angular momentum mixing hypothesis. J. Science (1965). Hurricane formation by convective stirring**
- **vorticity expulsion hypothesis, Gough and Lyndon-Bell, JFM (1968). Turbulence scrambles vortex lines and drives mean vorticity to zero and expels to the edge of the turbulent region.**



Convective overturning and potential vorticity (PV) conservation – a thought experiment

- **Consider initial rest state in a rotating system where $M = Ar^2$ and $\theta = \theta_0 + Bz$ where $A > 0$ and $B < 0$**

$$PV = \frac{1}{r} \frac{\partial M}{\partial r} \frac{\partial \theta}{\partial z} < 0$$

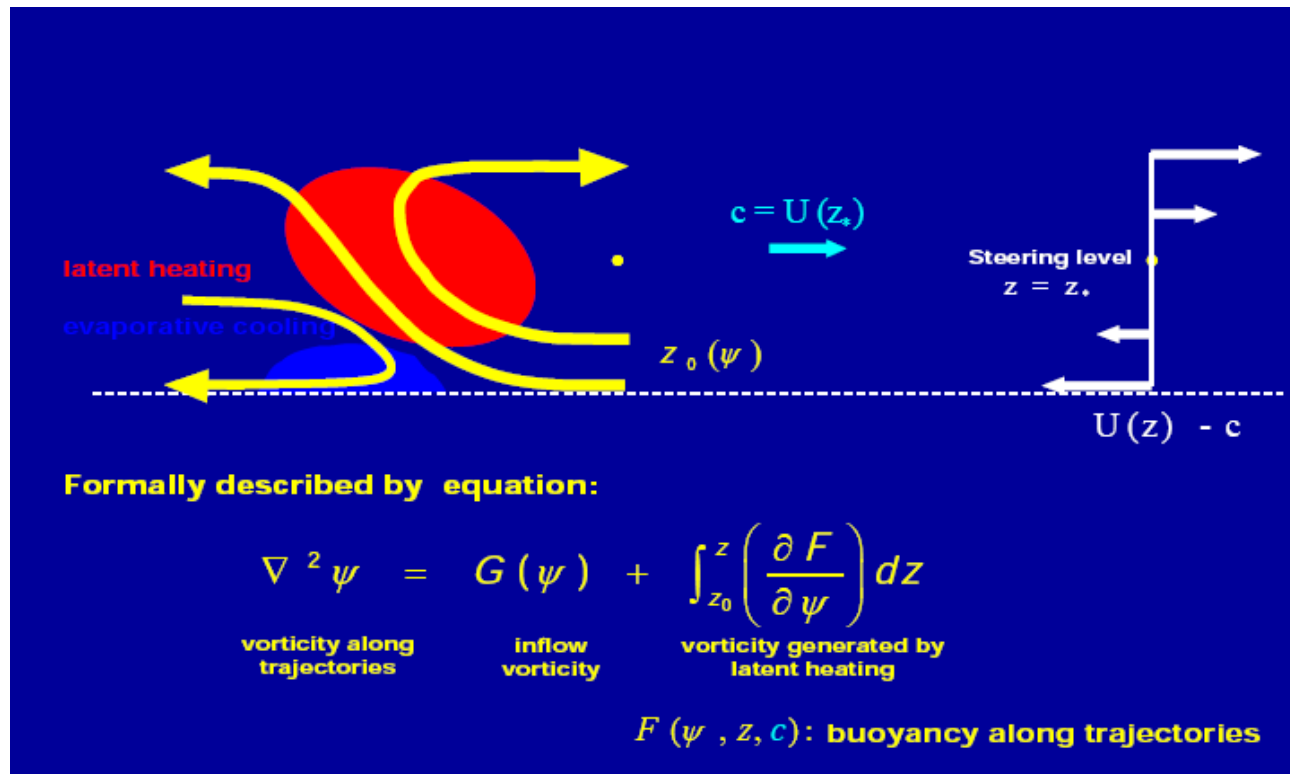
Overturning whilst conserving PV and global angular momentum would imply reversed radial gradient of M

e.g. $M = M_0 - A r^2 \longrightarrow$ **Vortex !**

Energetically impossible unless we exclude a cylinder of fluid at the origin

Cloud momentum transport

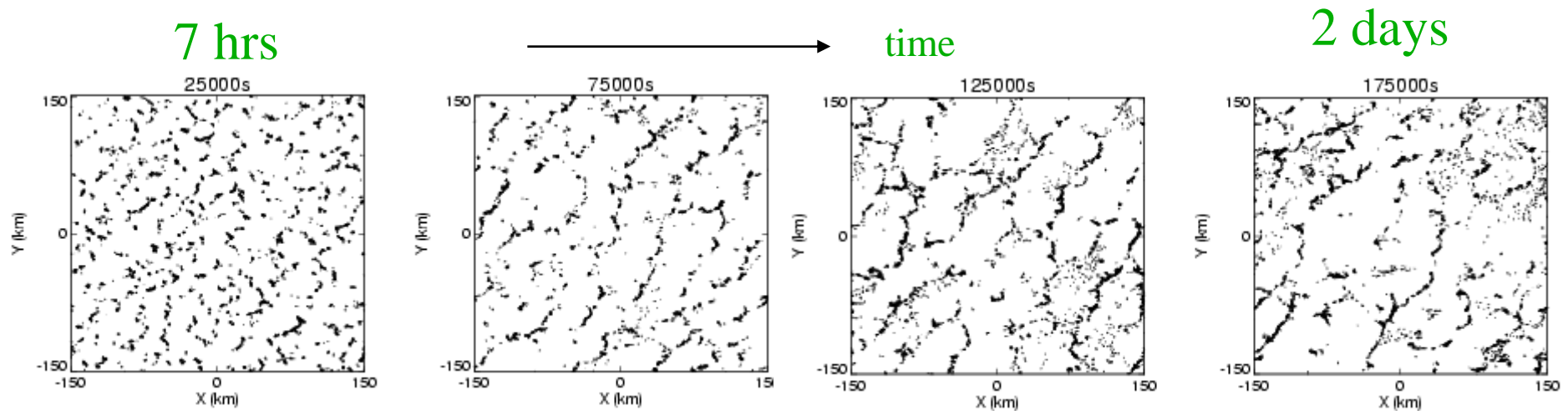
- Momentum not conserved on air parcels but vertical parcel exchange still causes downgradient transport
- Upgradient transport possible in squall line systems (Moncrieff and Green, 1972; Moncrieff, 1982 and 1991)



Upscale energy cascade

- Deep convective systems leave a mesoscale potential vorticity (PV) ‘footprint’
- PV anomalies have associated balanced flow fields
- Upscale energy transfer is caused by straining PV anomalies by large-scale flow

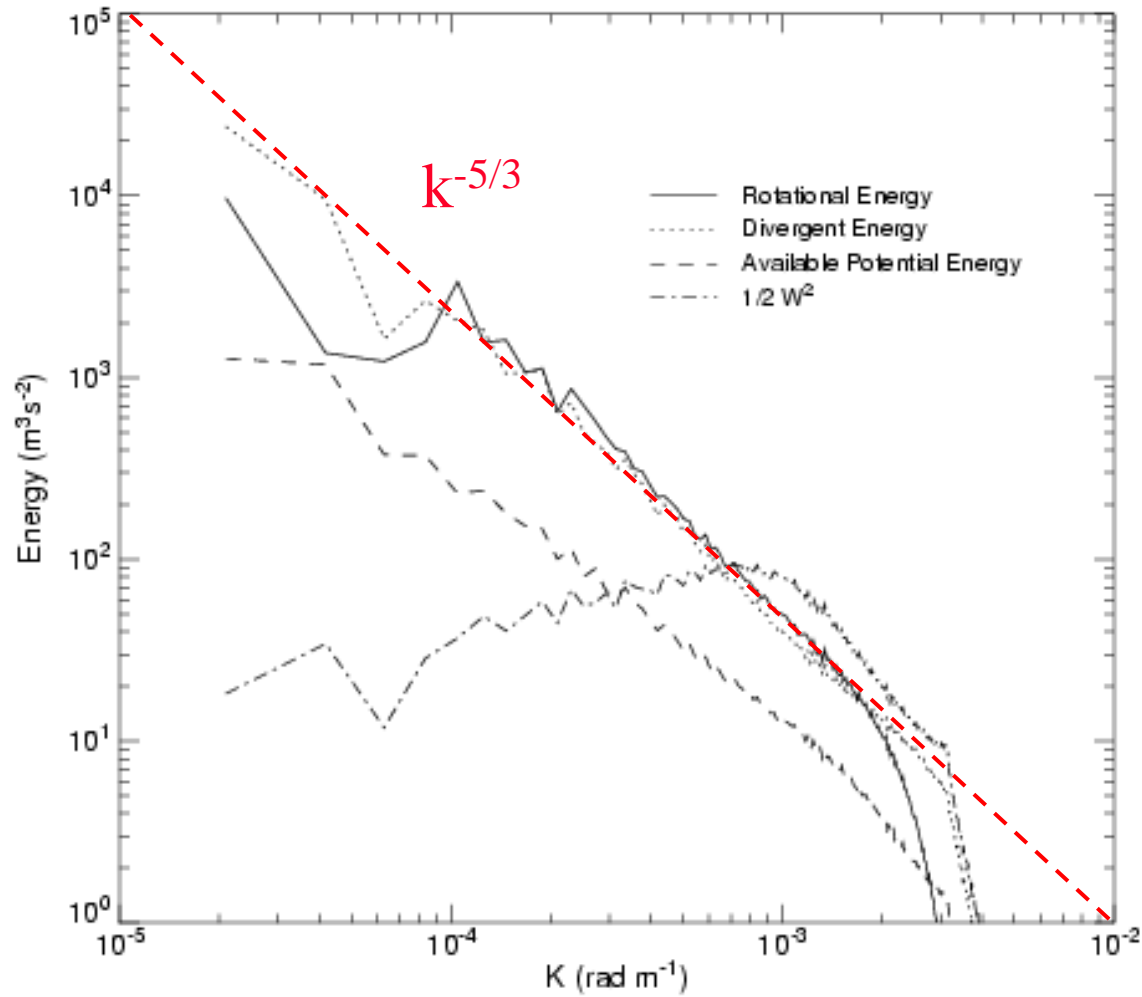
from Shutts and Gray (1999)



Deep convection / thunderstorm cloud - a view from space

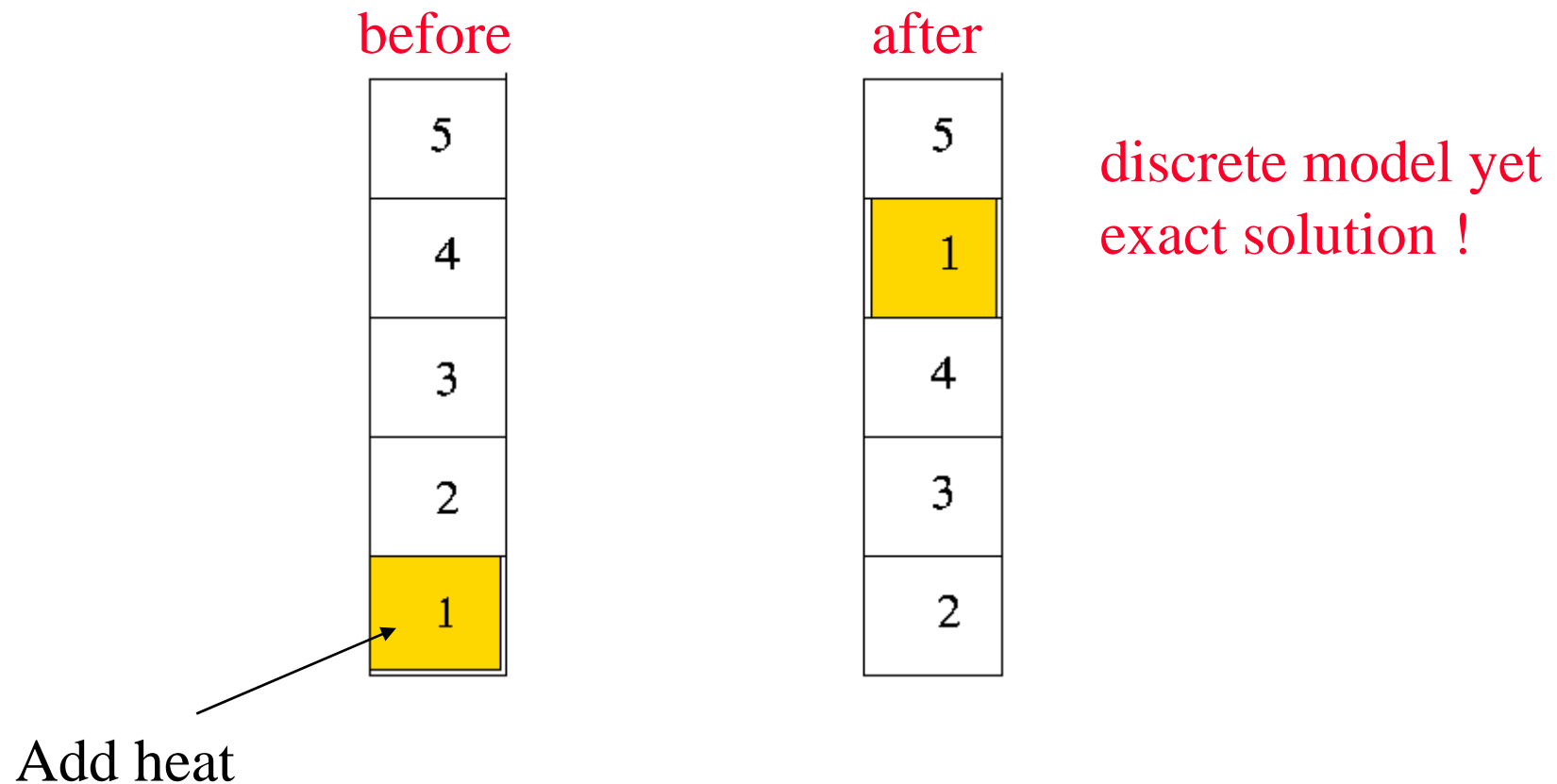


CRM simulation of upscale energy cascade (Vallis et al, 1997)



Parcel models of convective adjustment

- **Simplest model – constant volume and potential temperature lumps in a column**



2D parcel model in rotating system

- Air parcels conserved absolute momentum $M = \int v dx + \int u dy$ as well as potential temperature
- Inertial stability requires M increases monotonically with x
- Unlike 1D case we don't know parcel shapes *a priori* - only that they will be convex polygons
- the boundaries will be straight lines whose slope satisfies Margules's formula:

$$\frac{dz}{dx} = - \frac{f \theta_0 [M]}{[g \theta]}$$

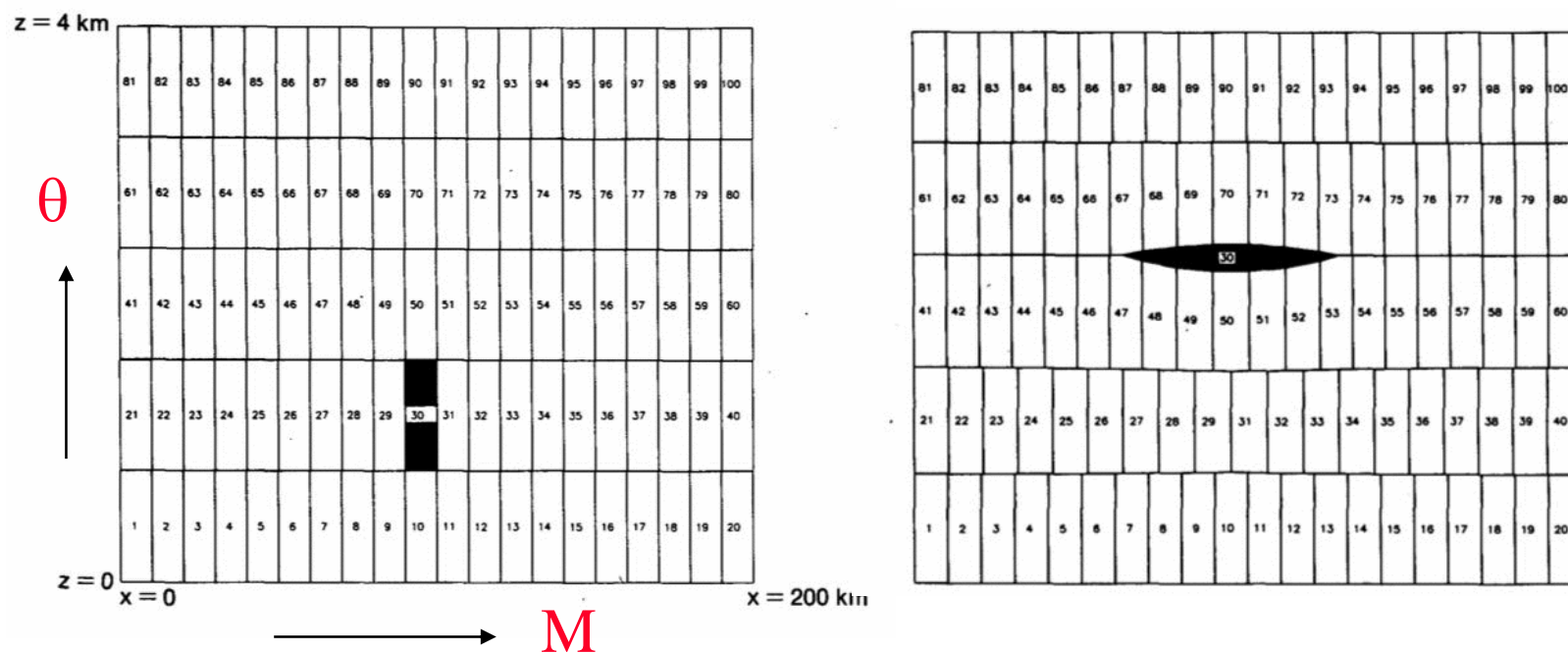
where $[]$ indicates the jump in value between parcels

2D parcel jump in rotating environment

uses Jim Purser's
element code

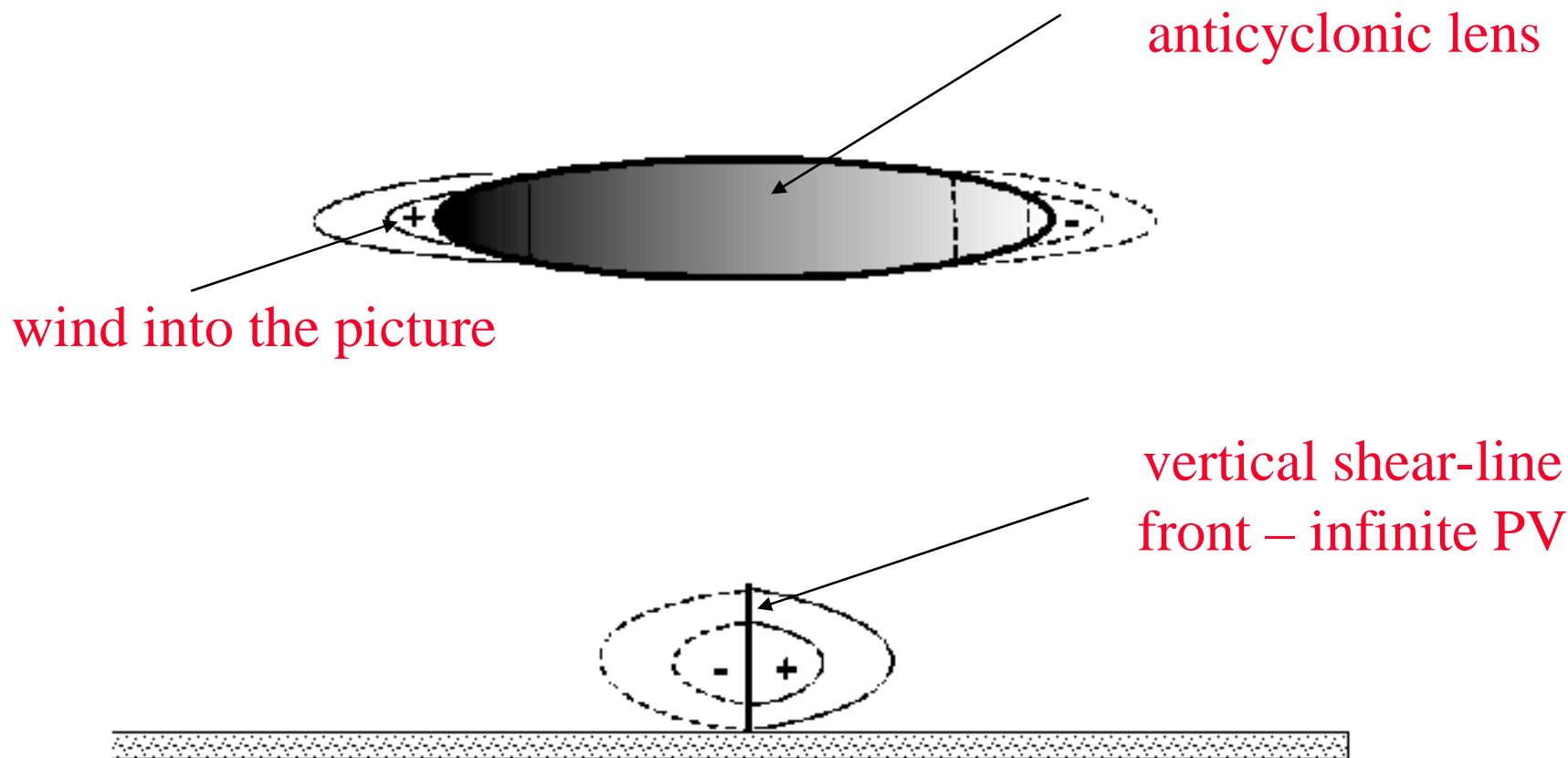
before

after



Note that rotation prevents the parcel from spreading into a thin block spanning the domain

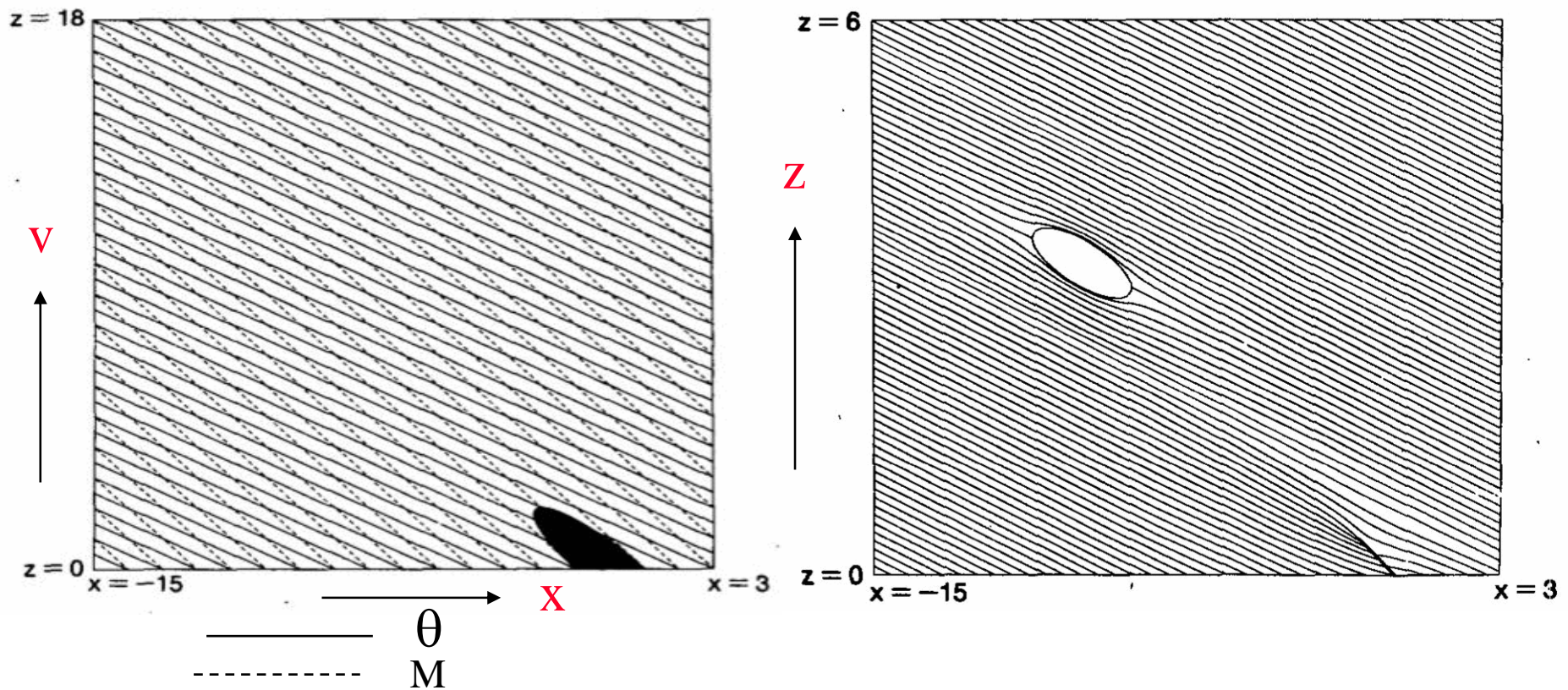
convective jump end-state



Slantwise convection parcel jumps

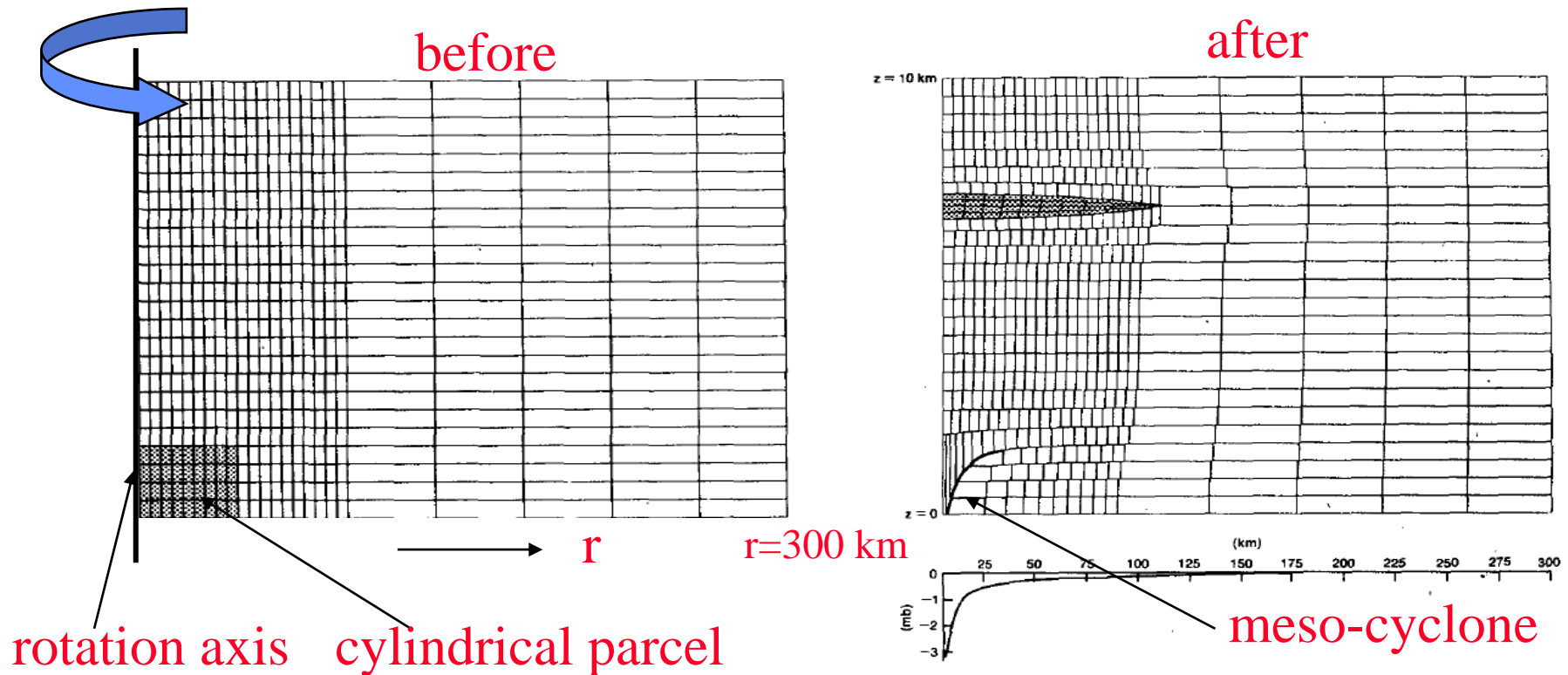
- 2D convection conserving M in an atmosphere initially with constant vertical wind shear in thermal wind balance
- Linear increase in θ with height.

from Shutts (1987)

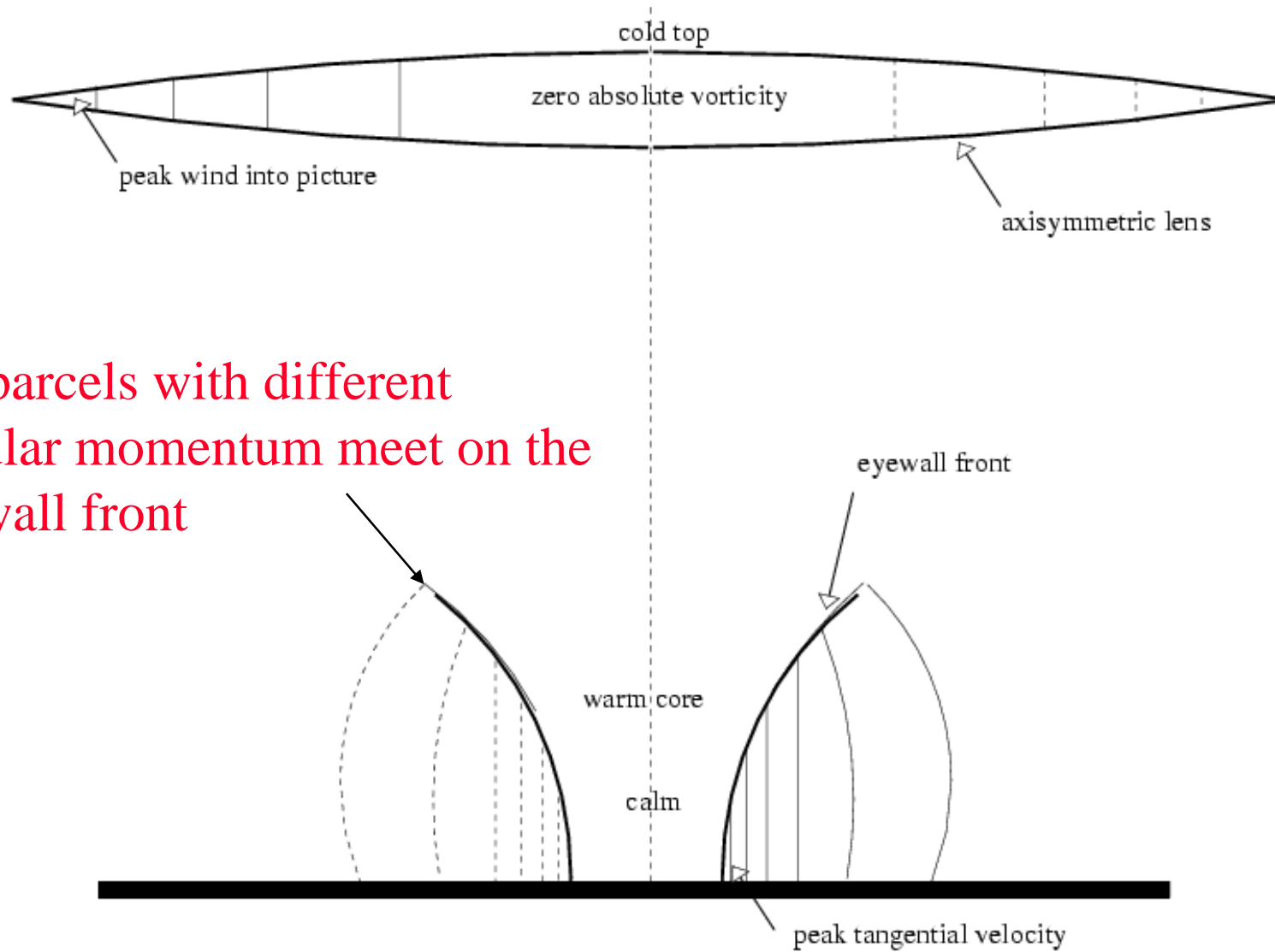


Axisymmetric parcel model

- Use **angular** momentum instead of absolute momentum
- Transform to 'bath plughole vortex' coordinates (makes parcel boundaries straight lines)
- Variable parcel sizes but conserve torus volumes



schematic picture of the end-state for a cylindrical convective parcel jump



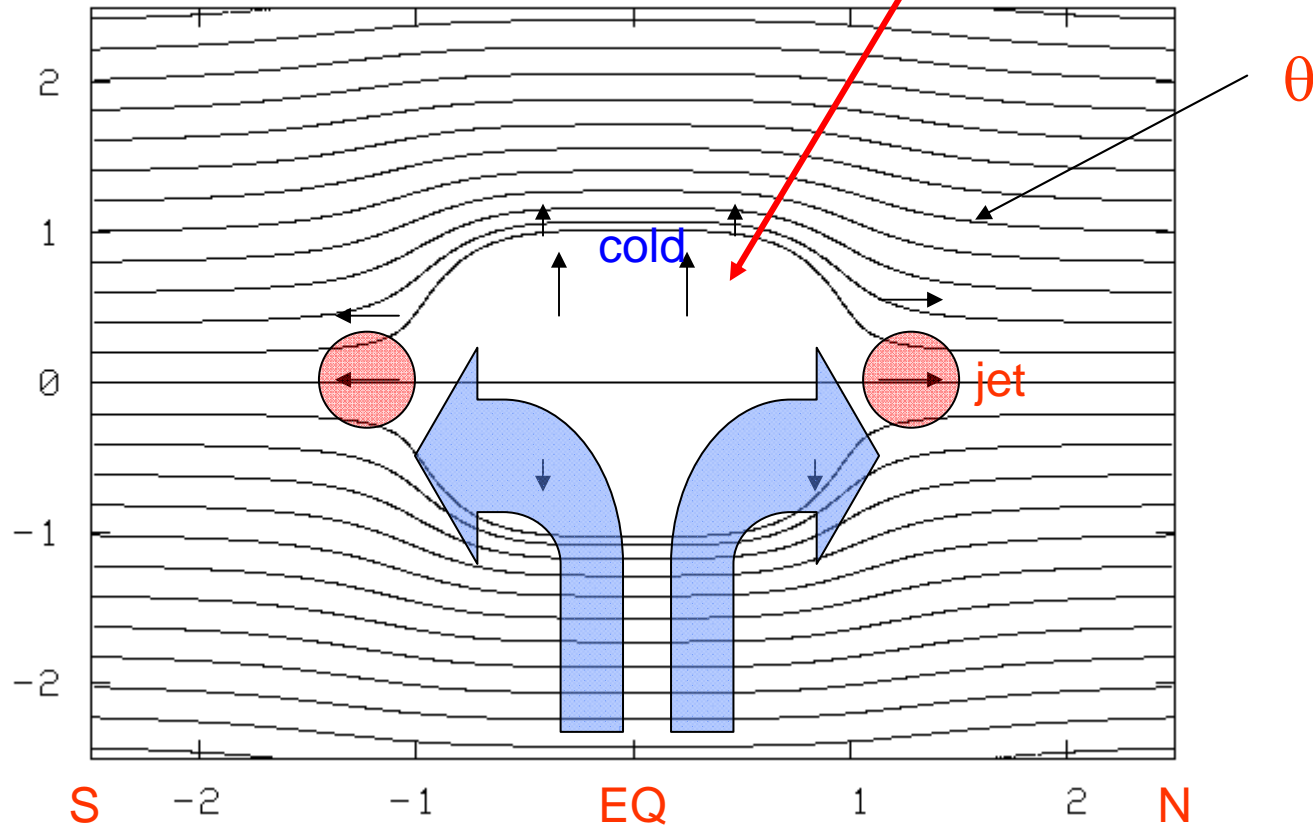
Air parcels with different angular momentum meet on the eyewall front

Convective mass flux – “pumping up the lens”

Homogeneous intrusion solution of Gill(1981) adapted for equatorial beta-plane i.e. $f = \beta y$

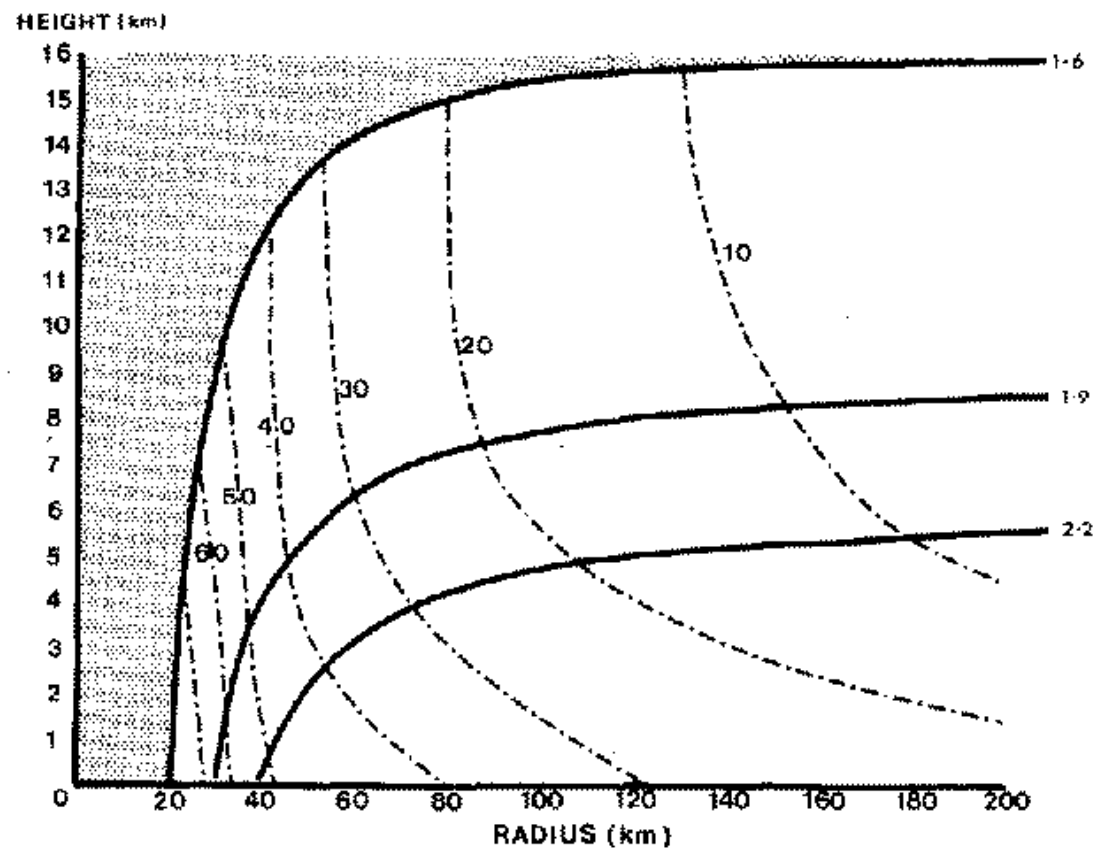
Zero PV region embedded in background linear meridional PV variation

Meridional-height section through the tropics



Hurricane structure – zero PV assumption

- **M and θ surfaces coincide**
- **Shape of surfaces fixed by $d\theta/dM$ and θ at $z=0$**

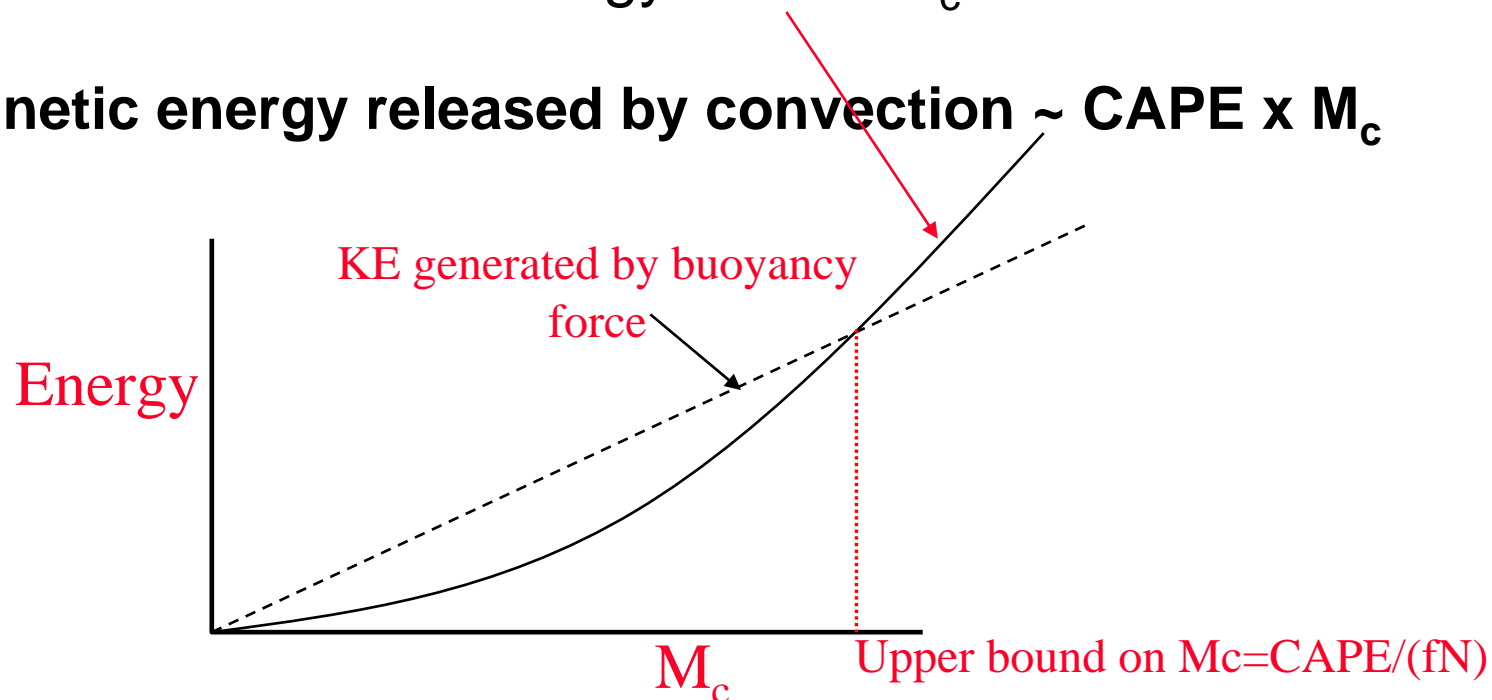


Mesoscale PV anomaly generation

- if lens has radius r , the thickness is $\sim (f/N)r$
- Velocity at lens rim $\sim fr$ and so $KE \sim (fr)^2 \times M_c$

where $M_c \sim (f/N)r^2$ is the mass convected
therefore the lens energy $E \sim fN M_c^2$

- kinetic energy released by convection $\sim CAPE \times M_c$



Convective length scales

max lens radius r_* $r_* = \frac{\sqrt{CAPE}}{f}$

CAPE=1000 J.kg⁻¹ and f=10⁻⁴ s⁻¹ → r_{*}=300 km

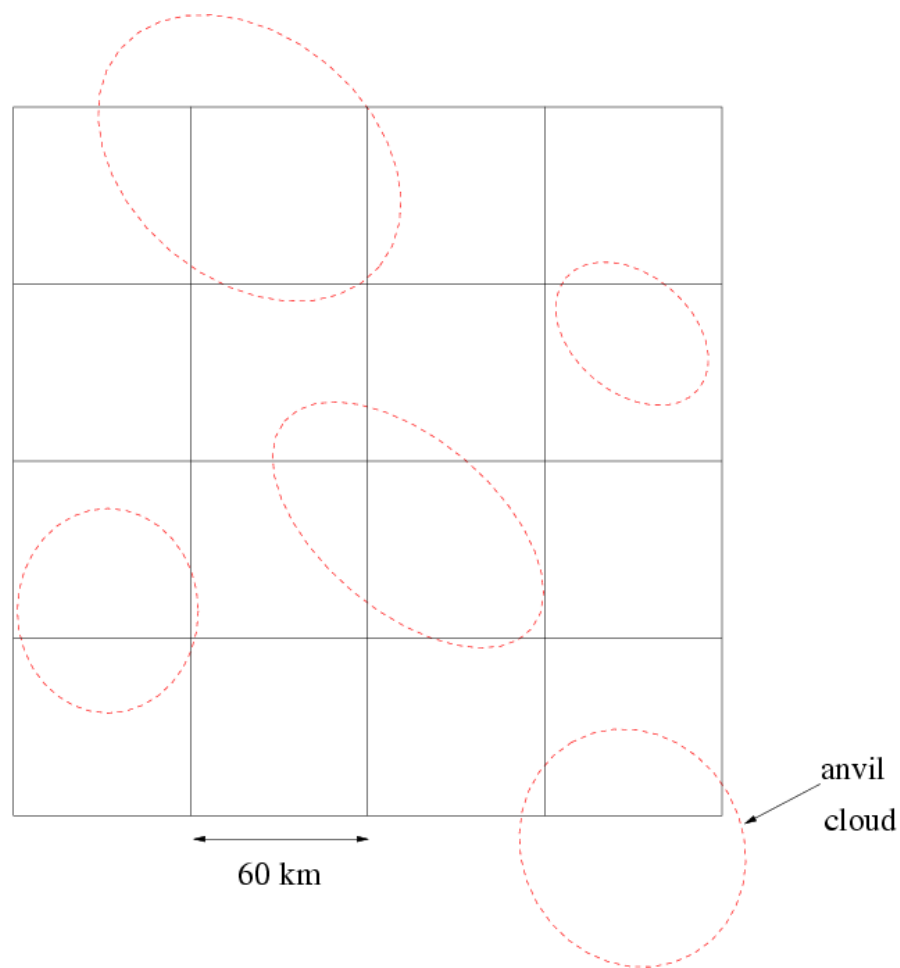
- Mesoscale convective system scale

Rossby radius of deformation (L_R) based on depth of convection (H_c) gives:

$$L_R = \frac{NH_c}{f} \sim 1000 \text{ km}$$

Convection parametrization issues

- **Low deep convective cloud density relative to 'gridpoint density**



Convection parametrization issues (continued)

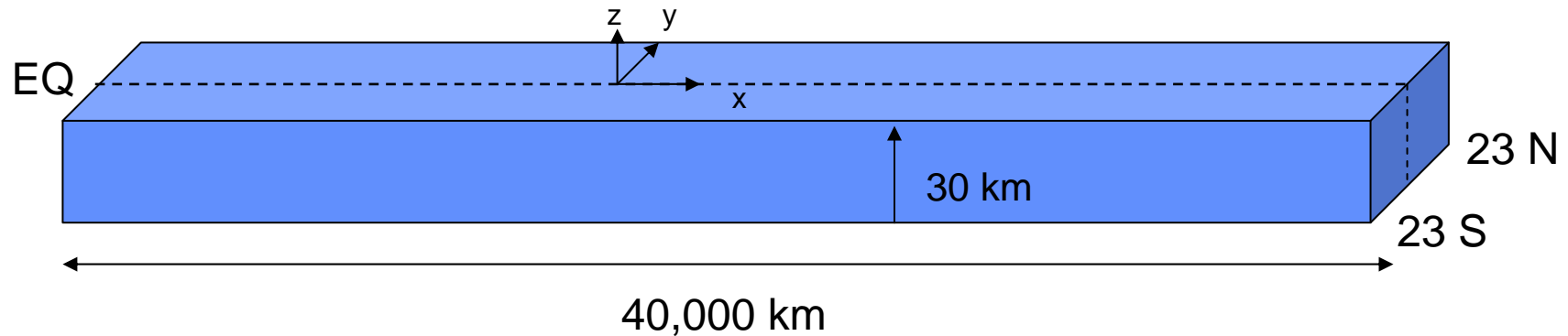
- **At the gridscale, are convective parametrization increments just noise ?**
- **Can current convective parametrization provide the correct upscale energy transports ?**

Use big-domain convection simulation to provide answers !

Big-domain simulation of tropical convection

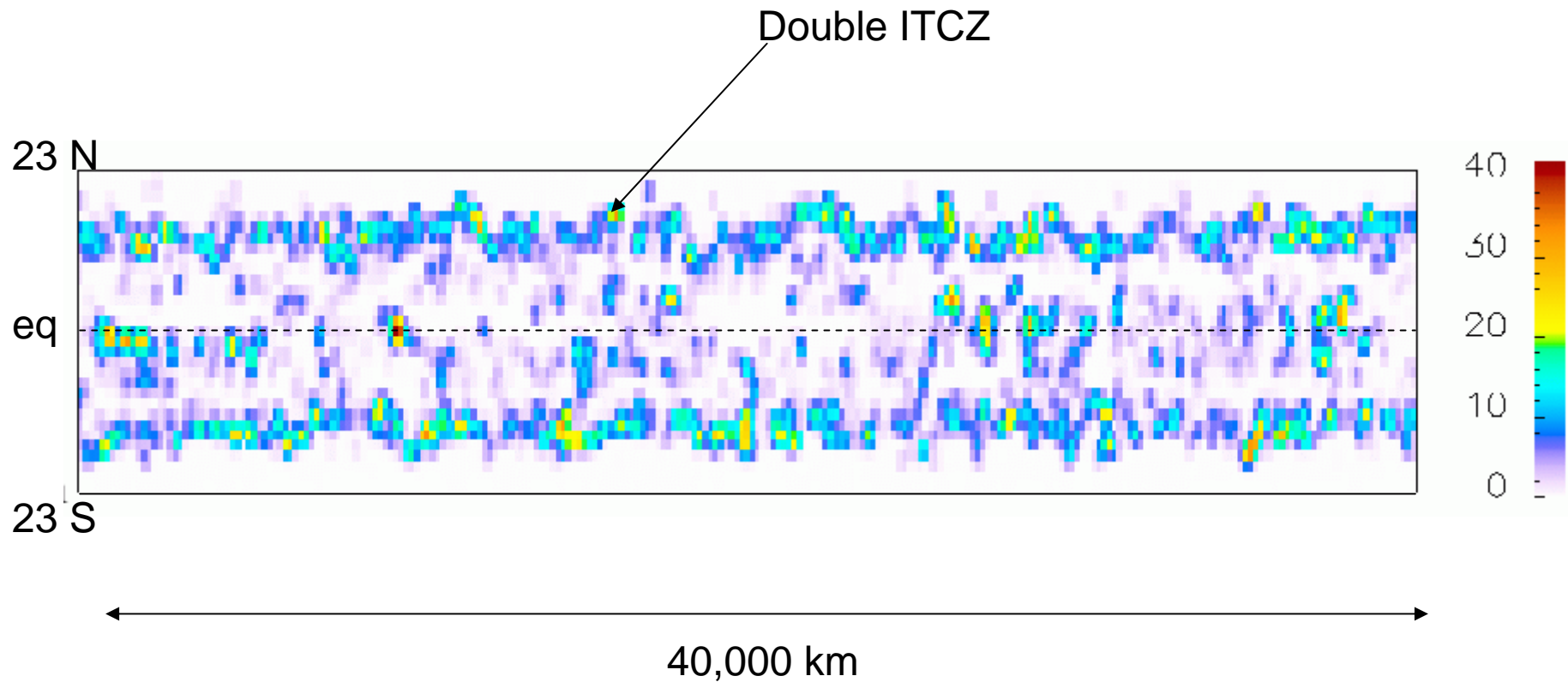
- attempt to simulate the interaction of deep tropical convection with large-scale flow with horizontal gridlengths > 1 km and domain sizes > 5000 km (in x & y)
- use $O(1$ km) resolution in x and $O(10$ km) in y
- run for at least 5 days but with short timestep (5 secs)
- coarse-grain fields *and* tendencies (i.e. source terms)
- compute PDFs, energy spectra, Fourier amp/phase plots

'circum-equatorial' model configuration

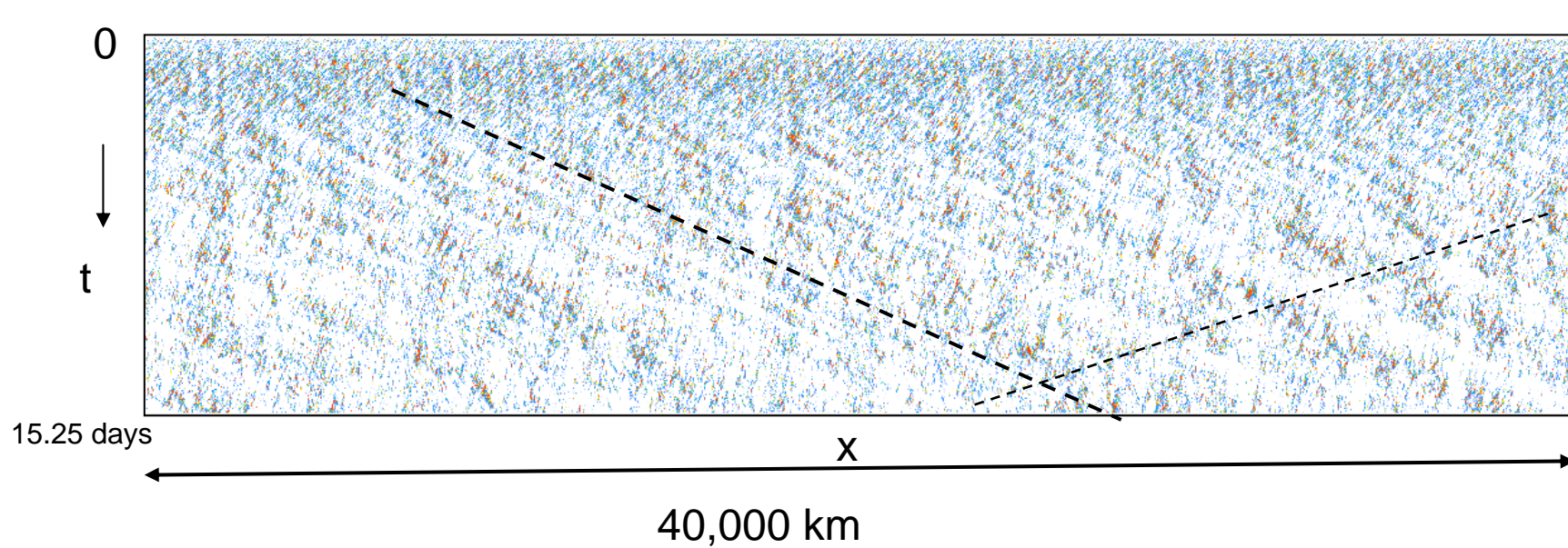


- $dx = 2.44$ km $dy = 40$ km 50 vertical levels
- 16384 x 128 x 50 gridpoints
- Coriolis parameter = βy
- impose 5 m/s easterly geostrophic wind
- fixed SST = $(28 - a y^2)$ degs C (a chosen so that N/S limits are 1.56 C cooler)
- no radiation, just imposed profile of cooling (-1.5 K/day up to 11 km)
- 3-phase cloud microphysics

Total rainfall over the 15.3 day CRM simulation

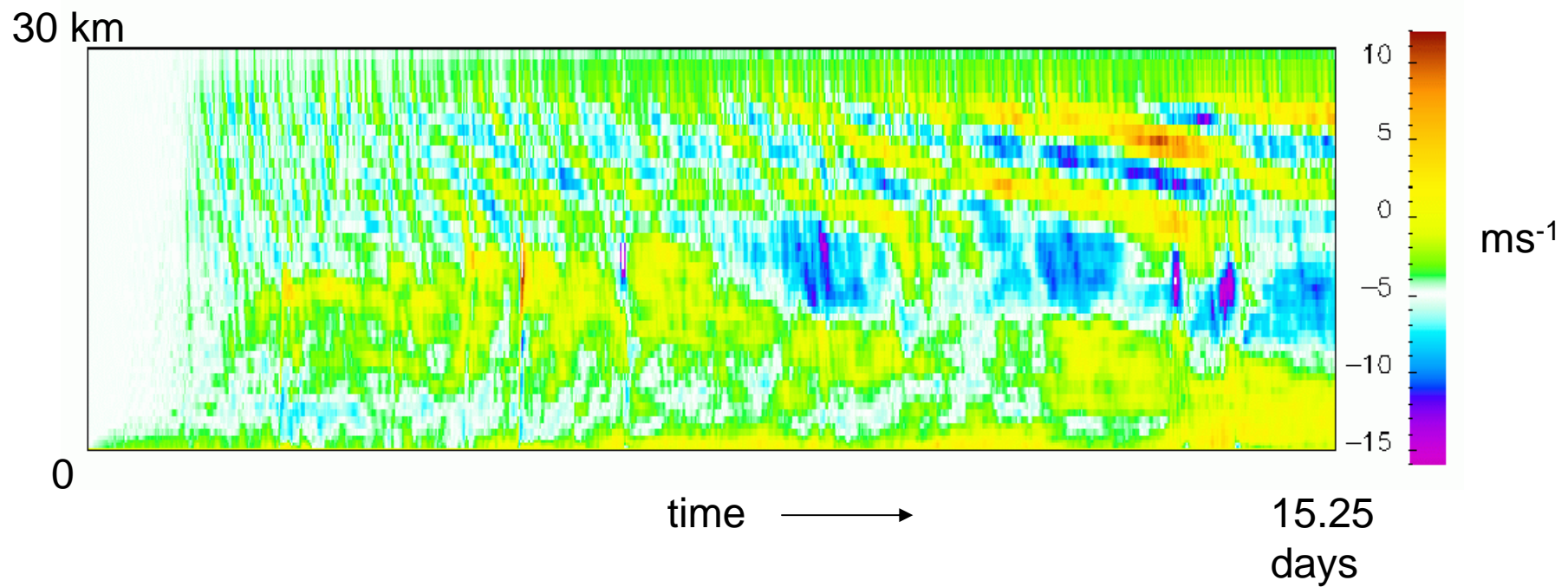


Hovmuller diagram of rainfall rate averaged over 10N-10S zone



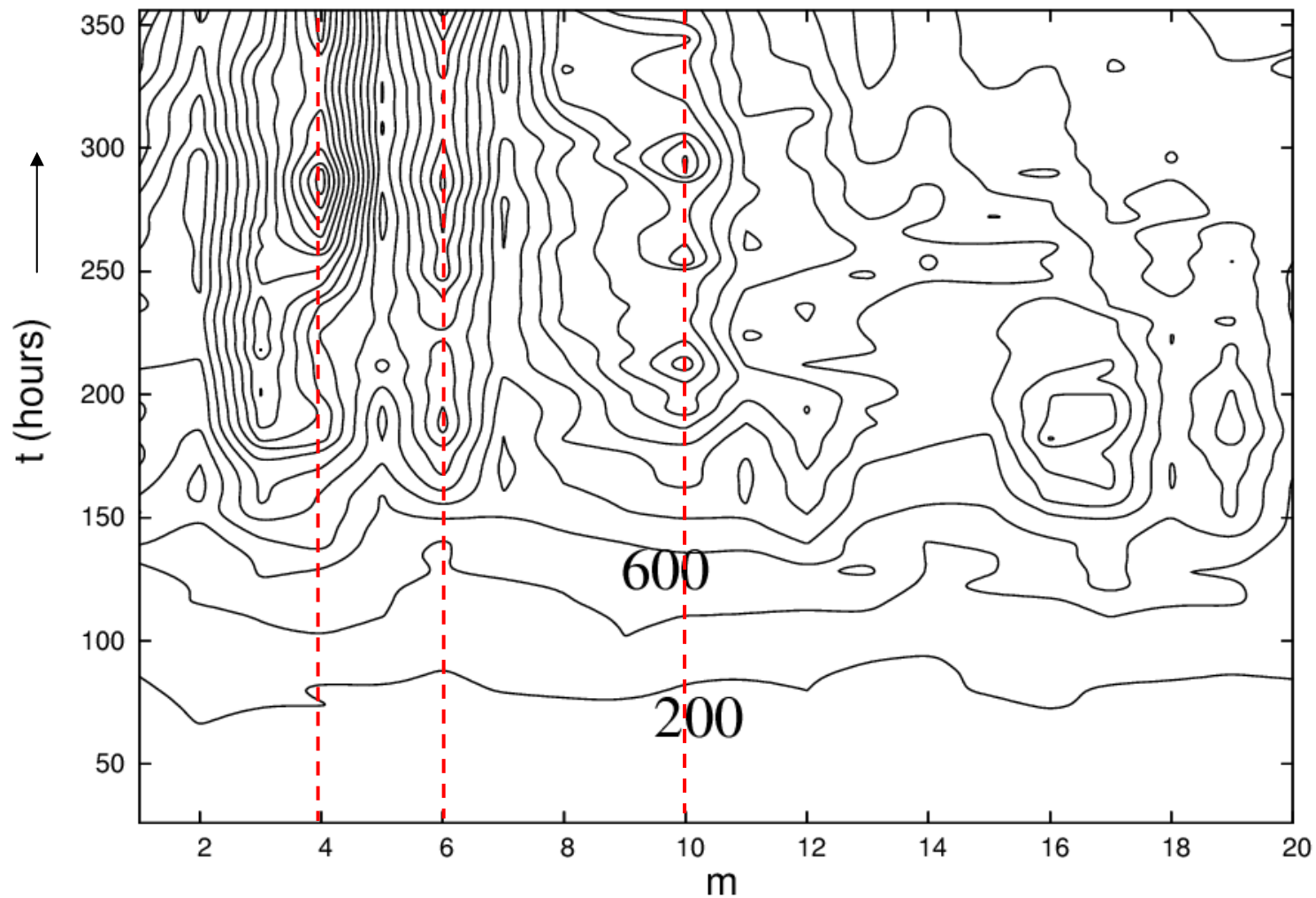
----- equivalent to 18 m/s propagation speed

Time-height section of zonal wind at a point on the equator



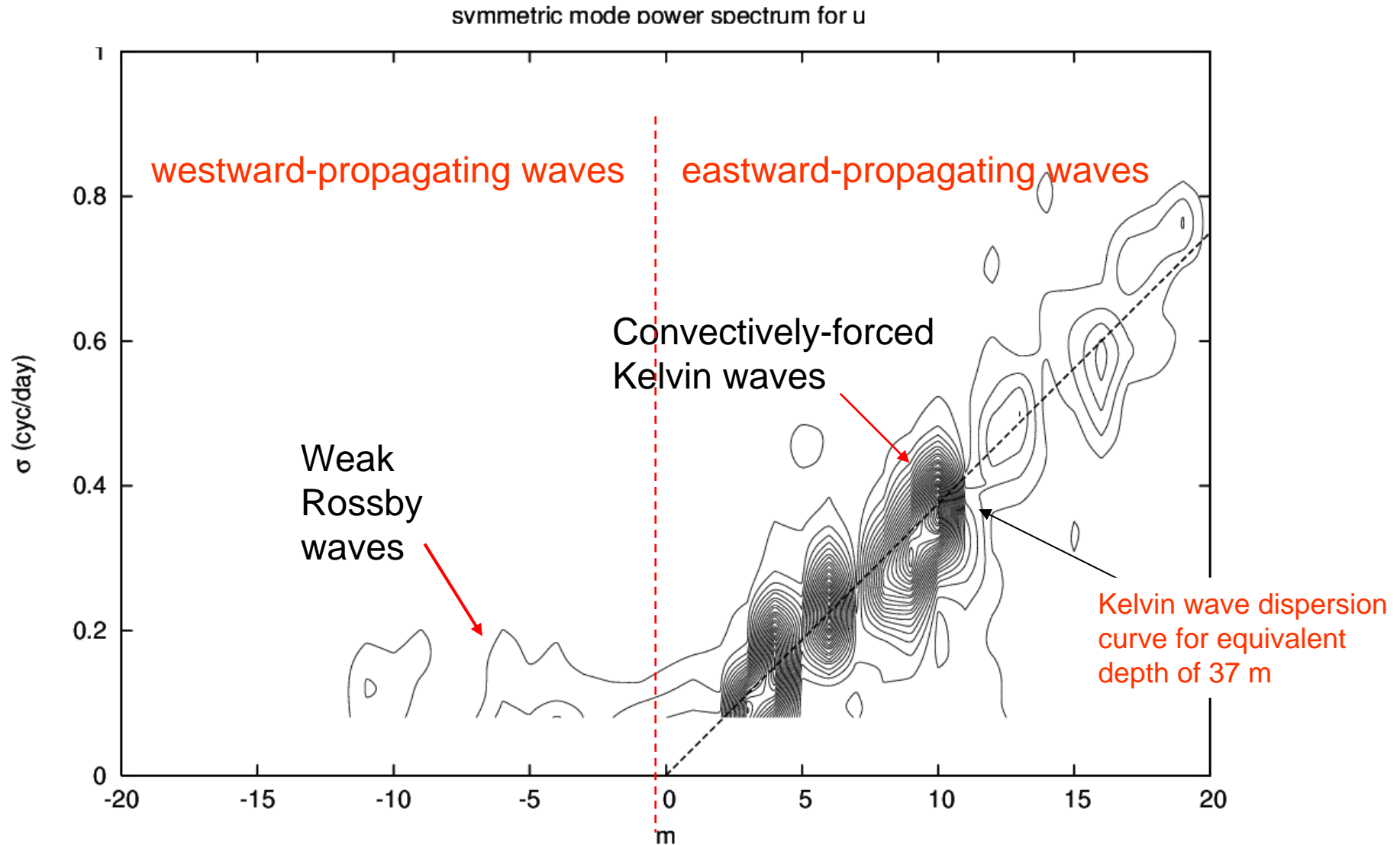
Growth of depth-integrated kinetic energy as a function of zonal wavenumber (m)

cont. int. 200 Jm^{-2}

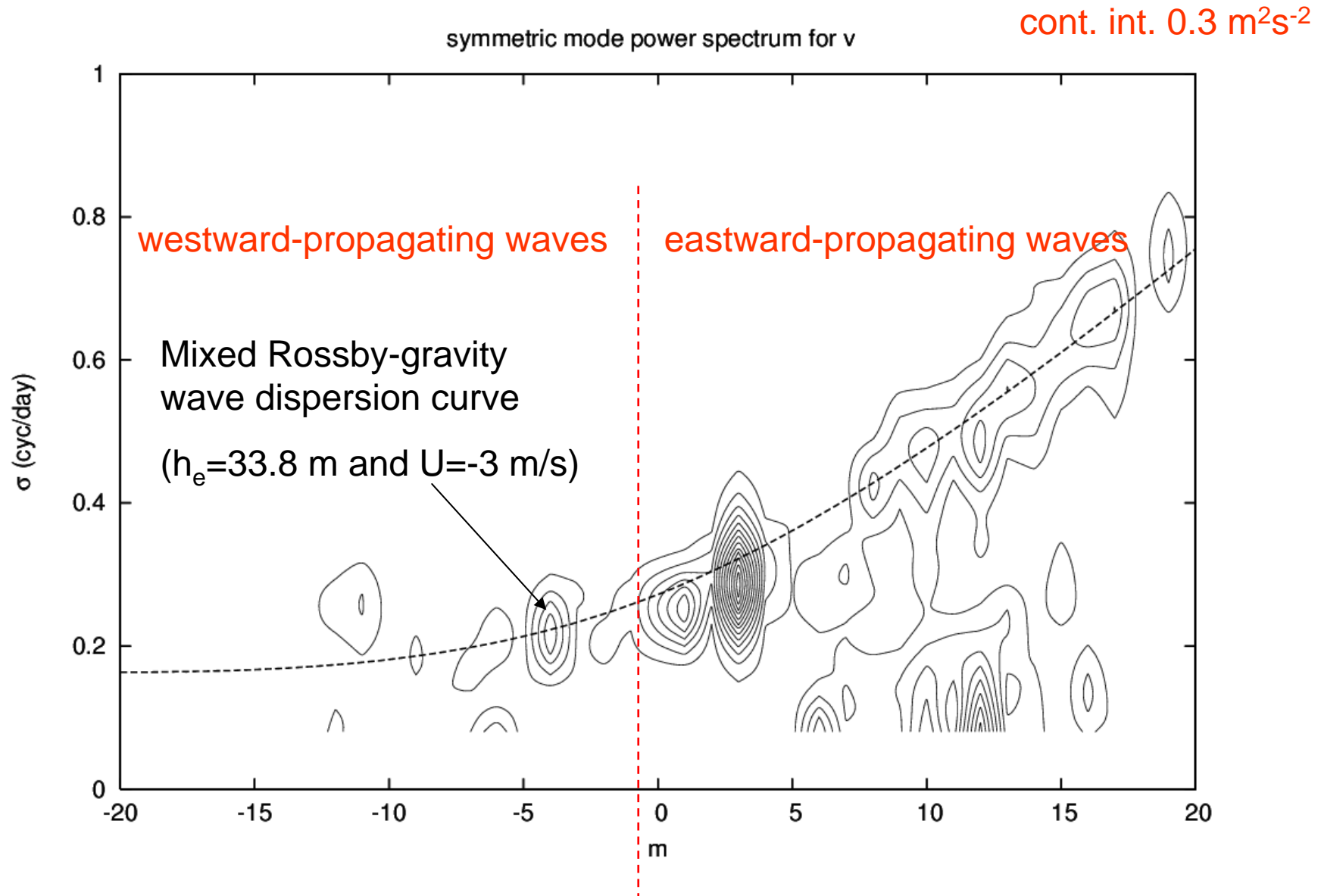


Symmetric contribution to the variance in u (height-mean)

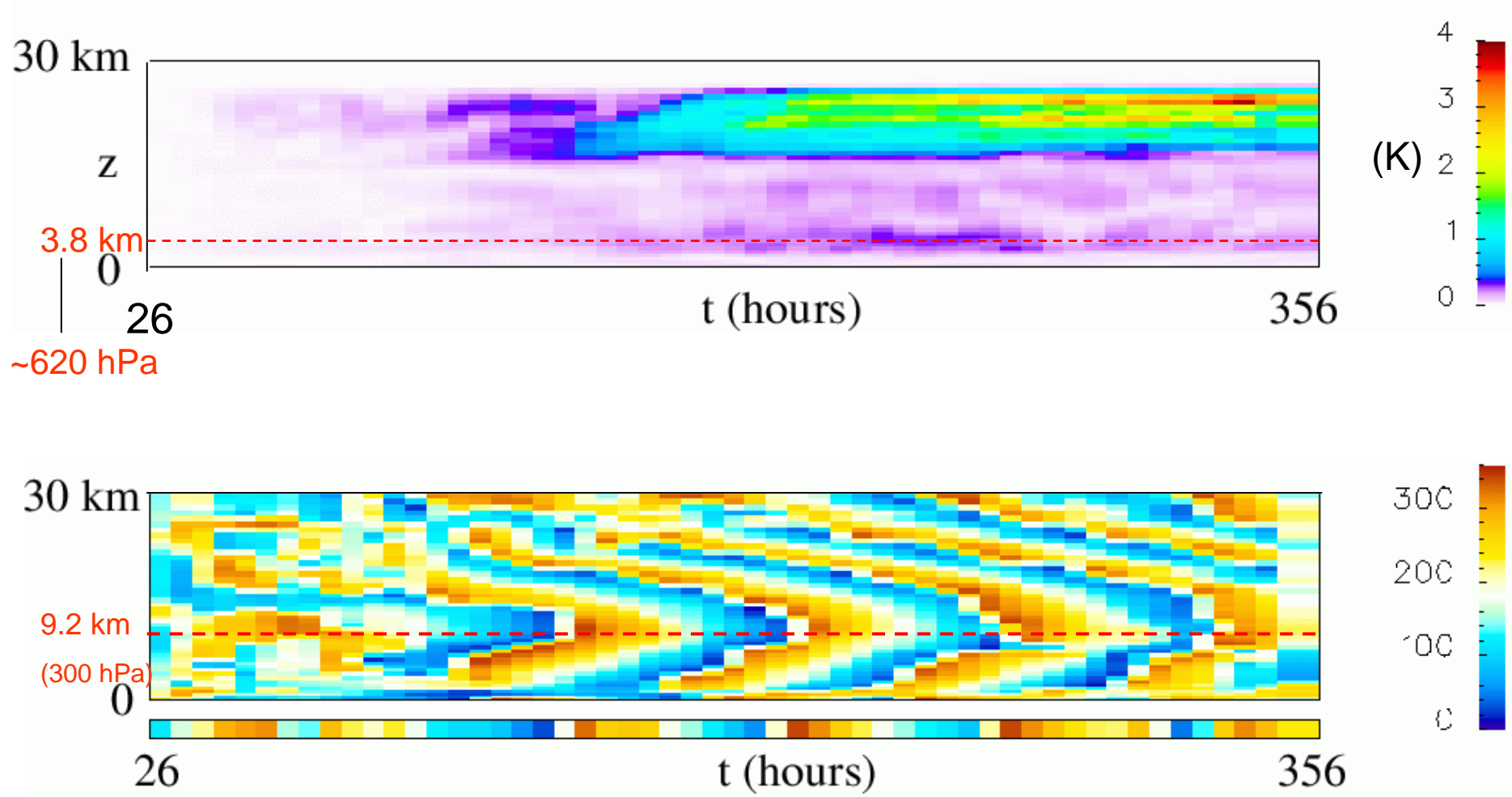
cont. int. $0.3 \text{ m}^2\text{s}^{-2}$



Symmetric contribution to the variance in v

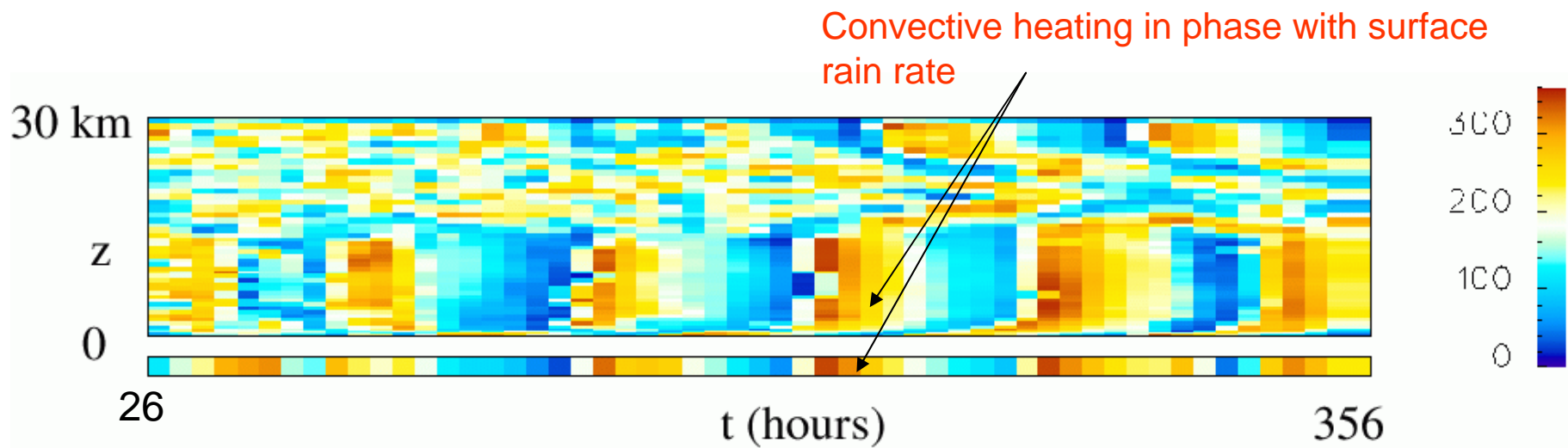
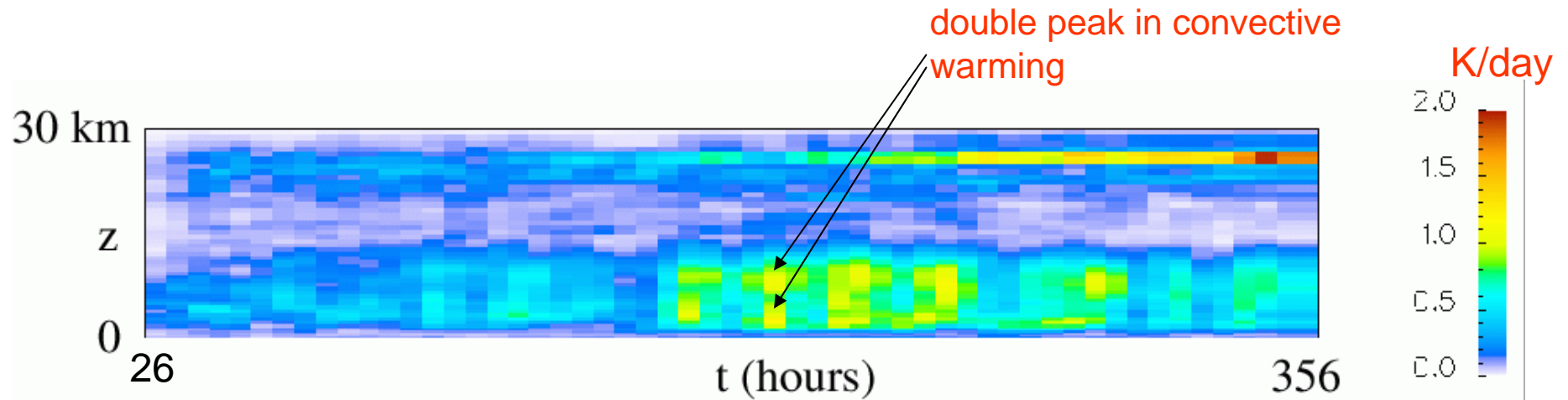


Time-height section of amplitude and phase of potential temperature perturbation. (m=10)

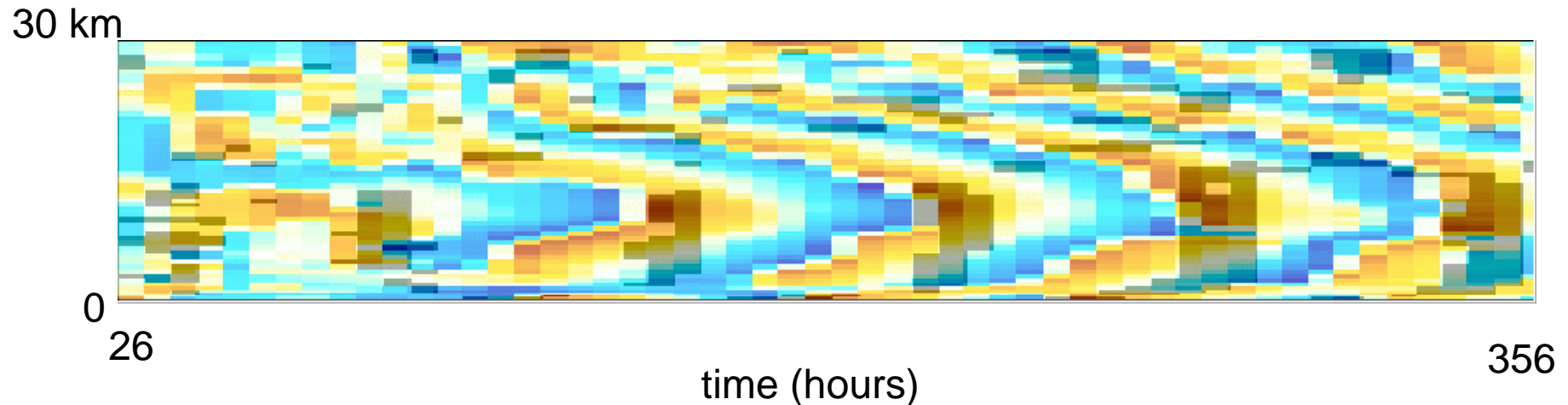


Time-height section of the amplitude and phase of Q

Zonal wavenumber 10

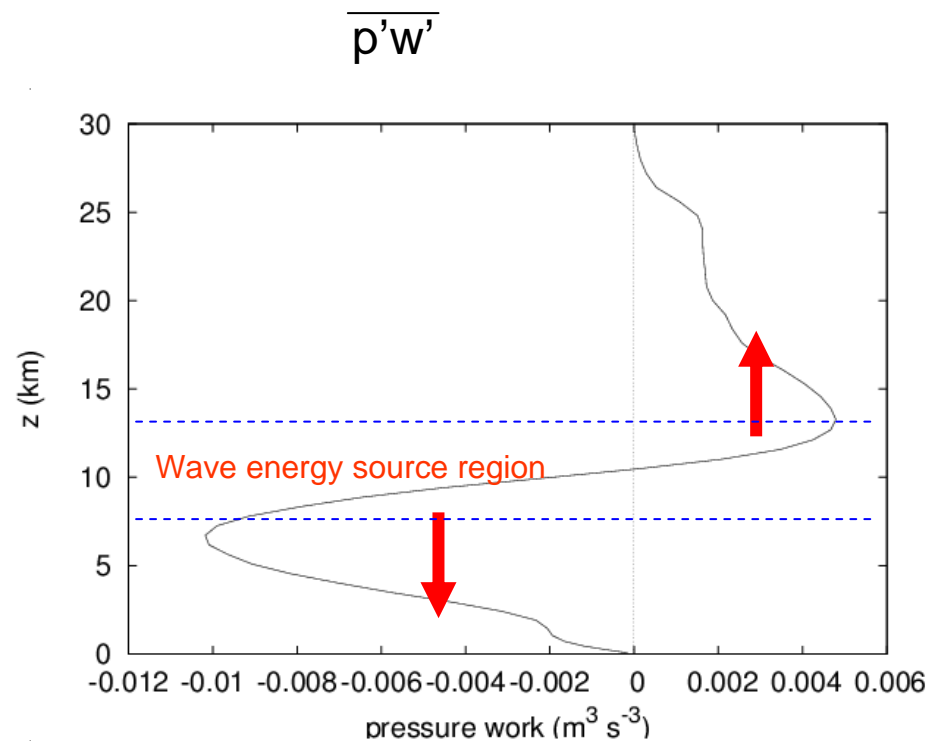
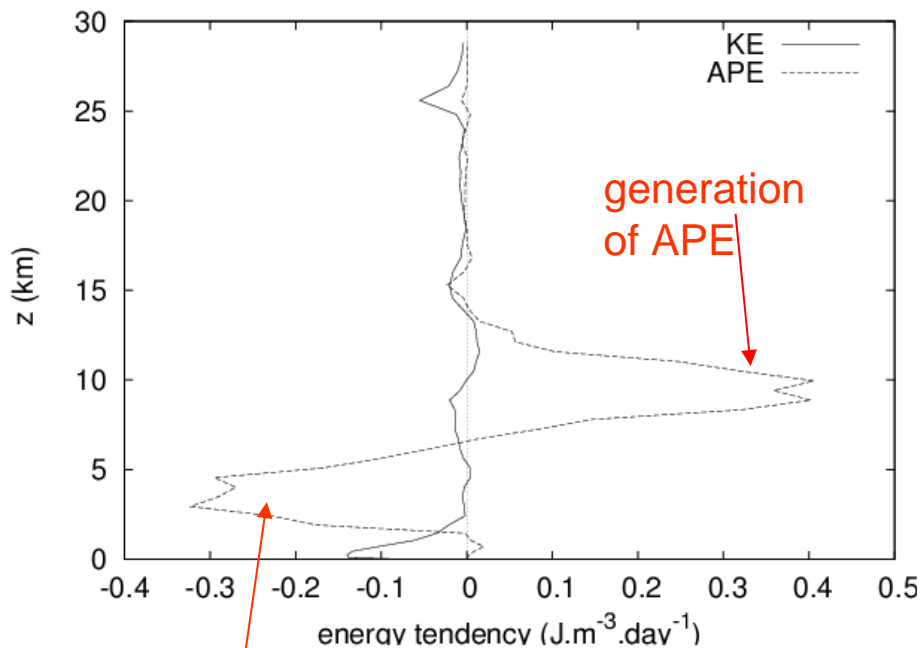


Composite of the time-height sections of wavenumber 10 phase for potential temperature perturbation and convective warming.



Think of red/orange as warm regions in $m=10$ wave
and dark shading represents convective warming

Vertical profiles of KE/APE production and pressure work at wavenumber 10



destruction of APE

Coarse-grain effective potential temperature tendency (\tilde{Q})

$$\frac{\partial \theta}{\partial t} = -\mathbf{V} \cdot \nabla \theta + Q$$

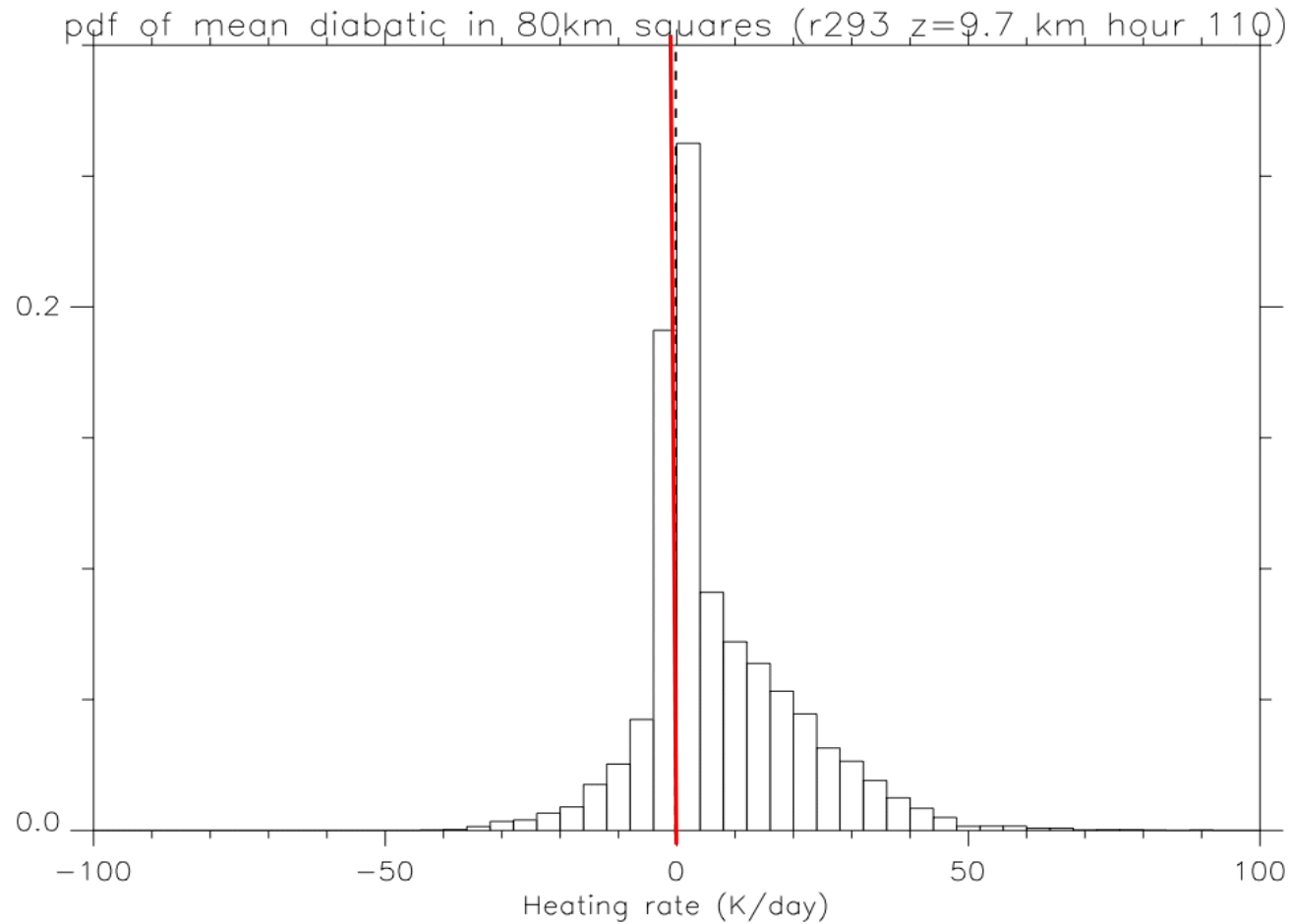
Let overbar denote average over a coarse grid box, then:

$$\frac{\partial \bar{\theta}}{\partial t} = -\overline{(\mathbf{V} \cdot \nabla \theta)} + \bar{Q}$$

$$\frac{\partial \bar{\theta}}{\partial t} + \bar{\mathbf{V}} \cdot \nabla \bar{\theta} = \underbrace{\bar{\mathbf{V}} \cdot \nabla \bar{\theta} - \overline{(\mathbf{V} \cdot \nabla \theta)}} + \bar{Q} = \tilde{Q}$$

Parametrized + resolved heating

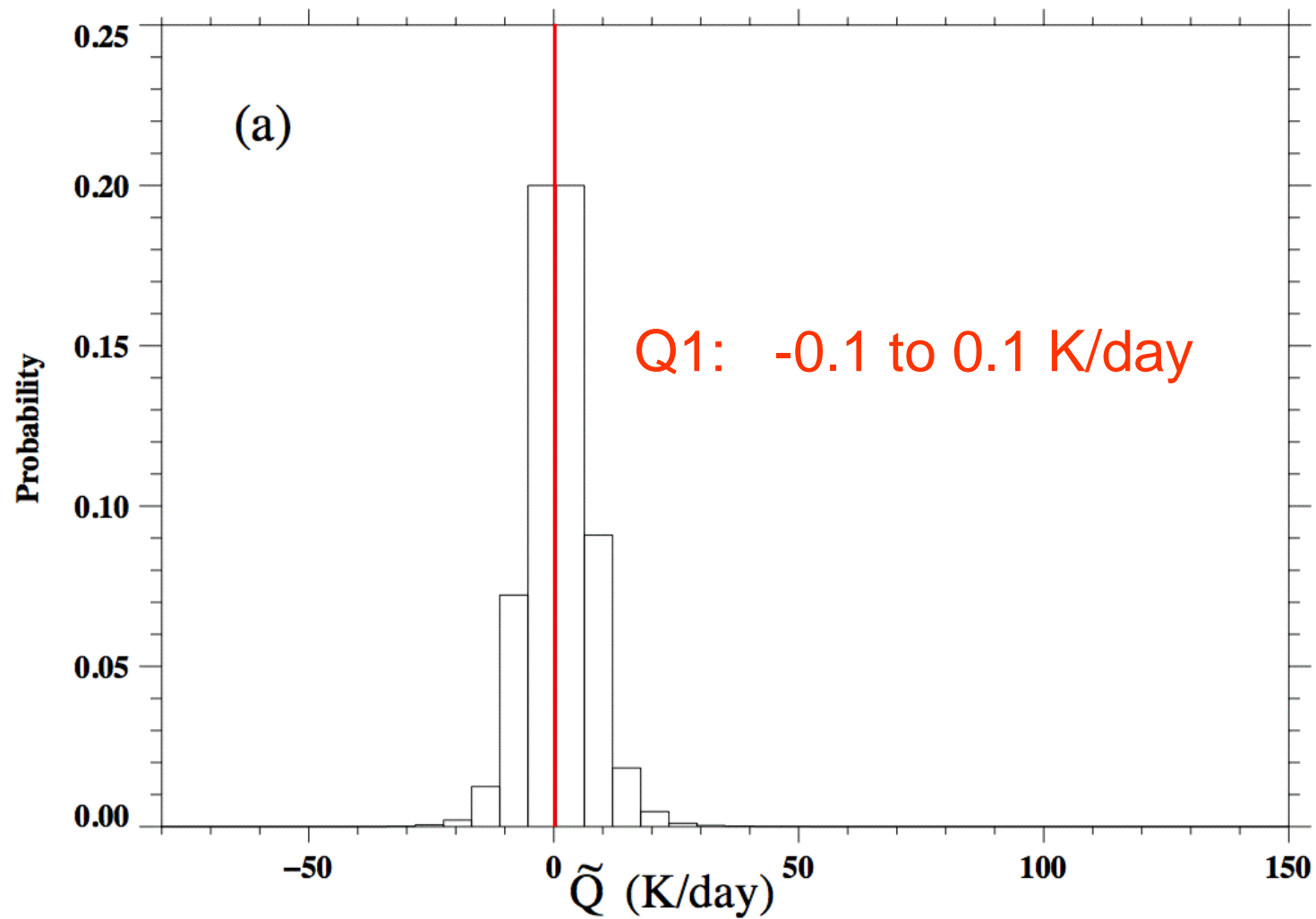
Histogram of diabatic heating (Q) coarse-grained to an 80 km grid at $z=9.4$ km

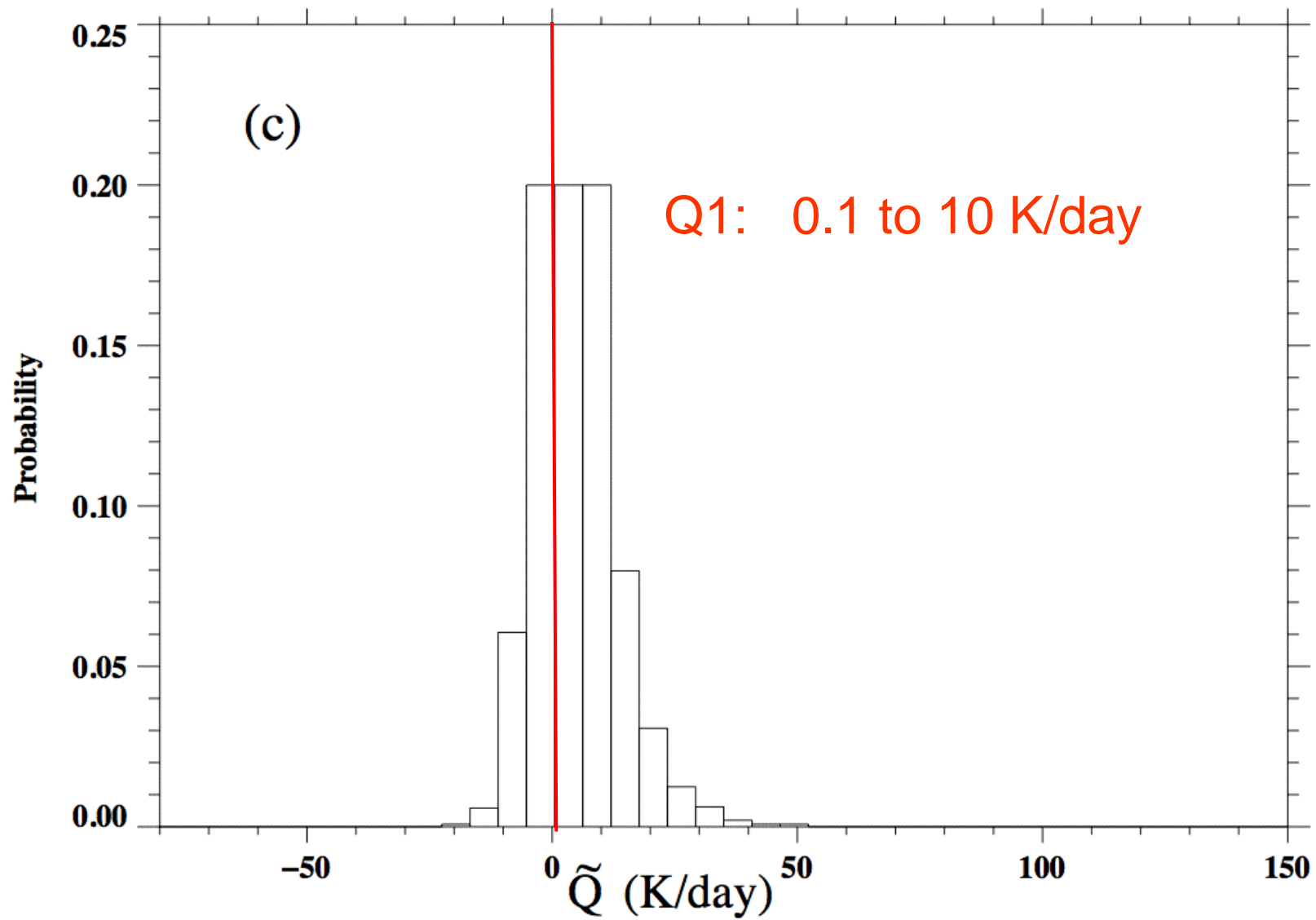


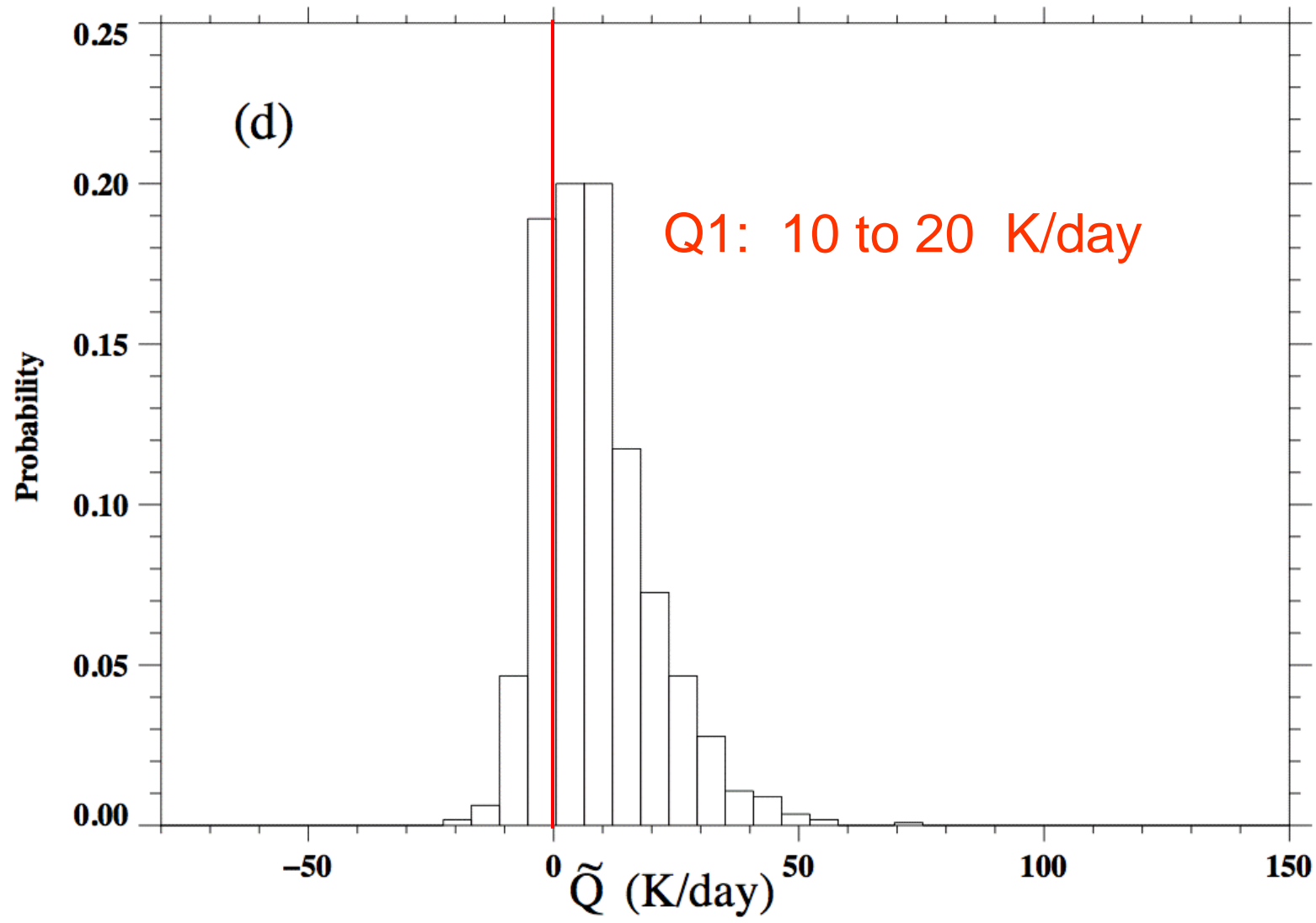
PDFs conditioned on convective parametrization temperature tendencies (Q_1)

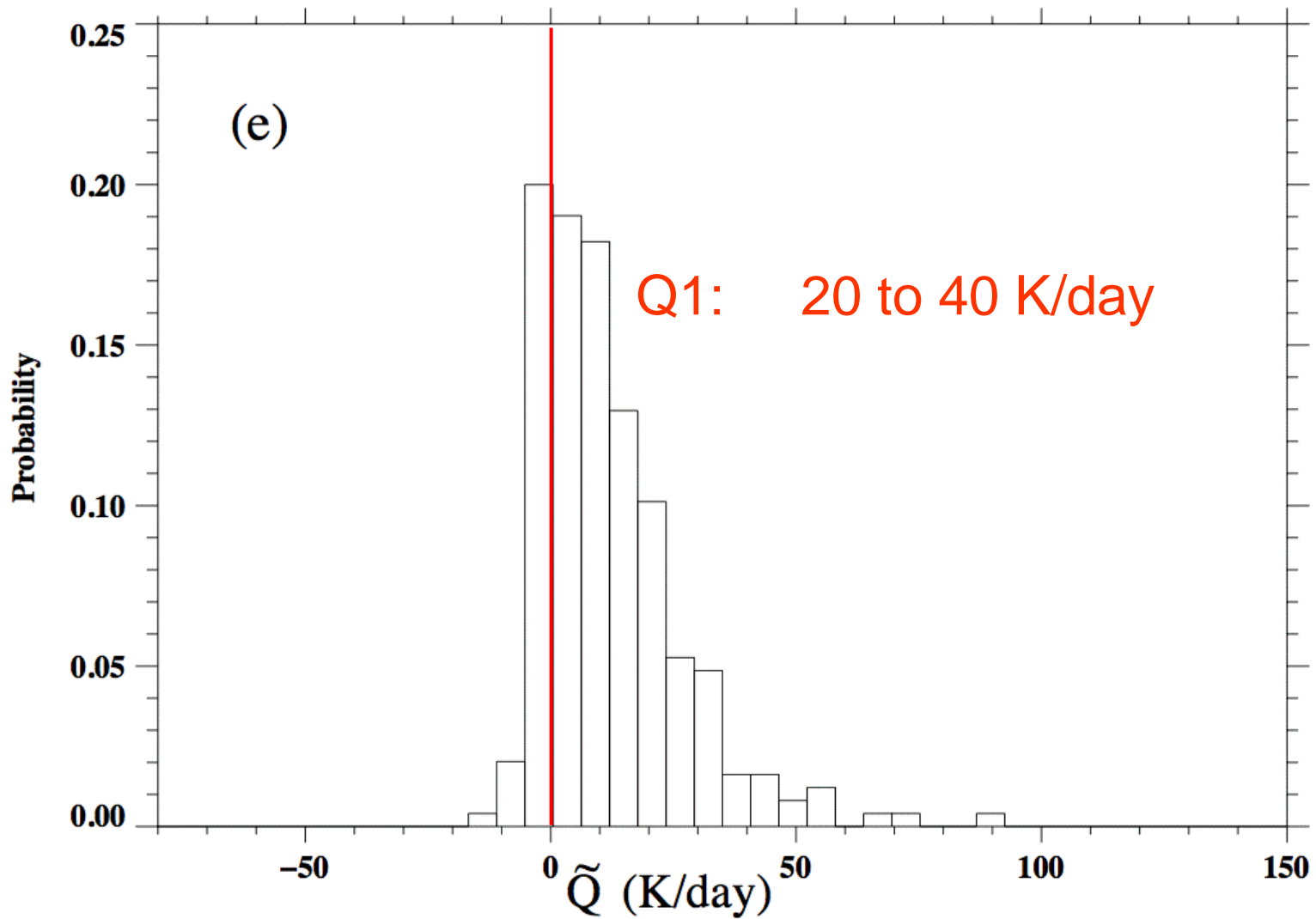
- take the coarse-grained CRM fields and feed them into a convective parametrization scheme (Bechtold et al, 2001) → Q_1 - the convective warming rate
- at any model level, bin the diabatic tendency \tilde{Q} according to different ranges of Q_1
- See how the variance of \tilde{Q} depends on Q_1
- Use knowledge of variance dependence to calibrate 'stochastic physics' schemes based on **multiplicative noise**

Pdfs of Q conditioned on different ranges of Q1

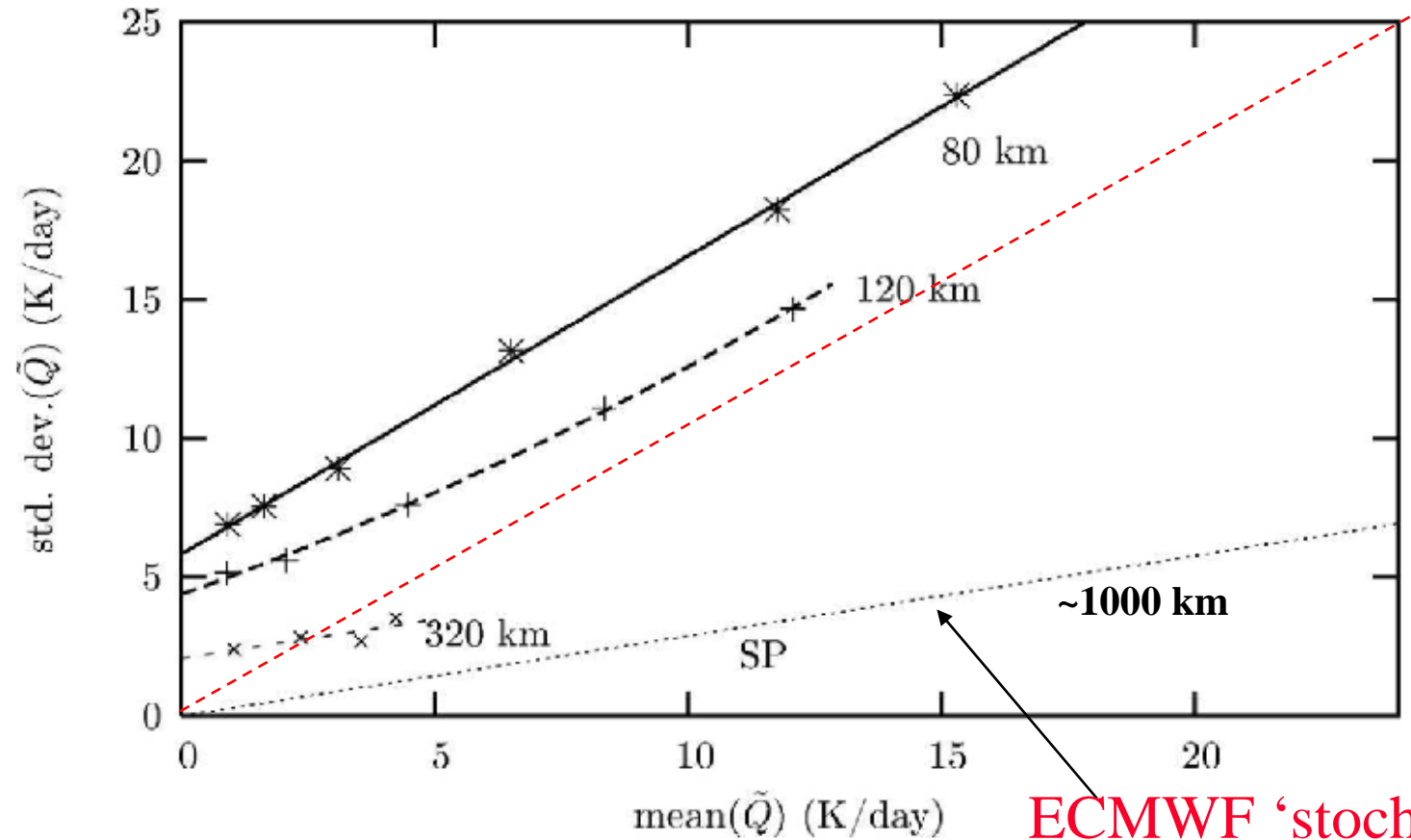








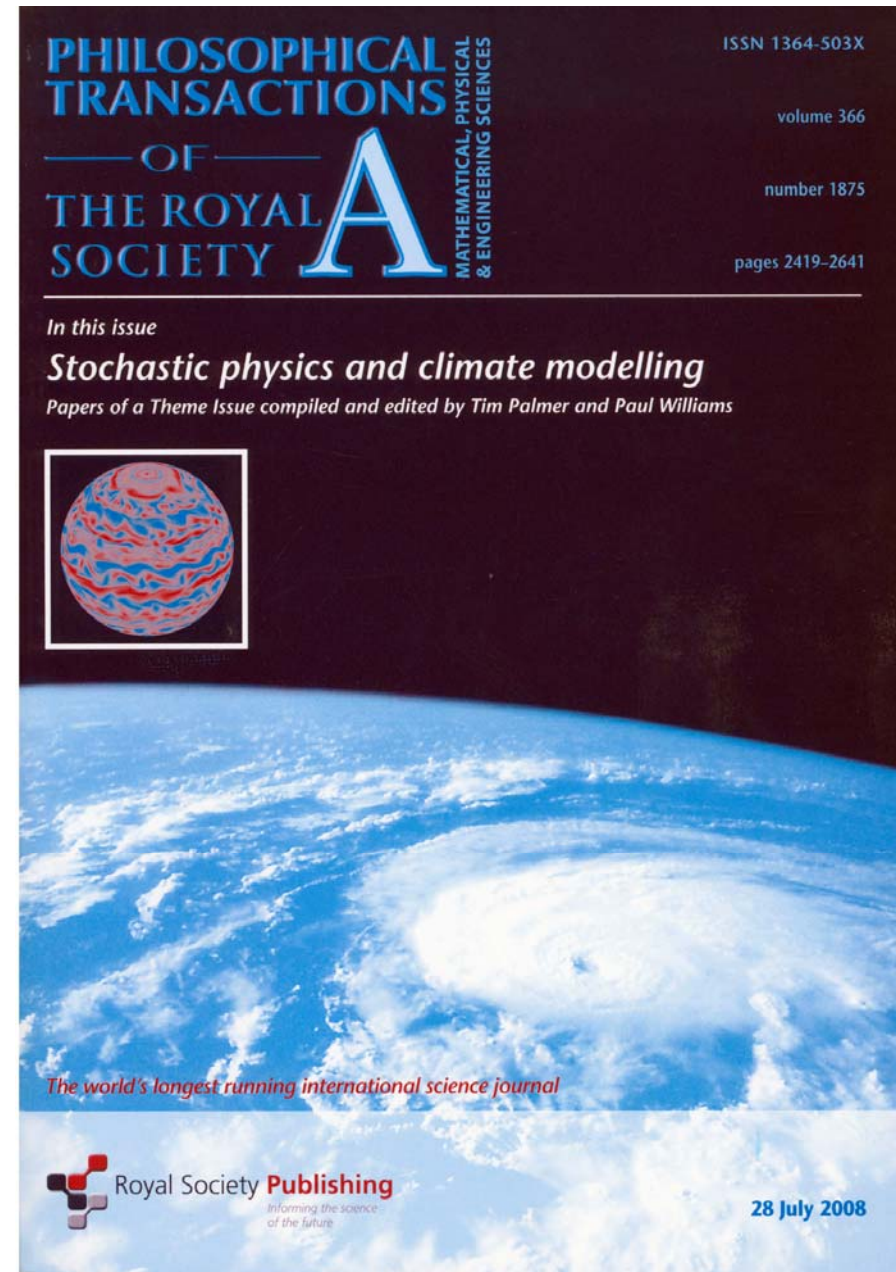
Variance of coarse-grained diabatic tendency



ECMWF 'stochastic physics scheme'

Phil. Trans paper

Volume 366, Number 1875
/ July 28, 2008 Theme Issue:
“Stochastic physics and climate
modelling” compiled by Tim
Palmer and Paul Williams



Stochastic convection parametrization

- **Buizza et al (1999)** -- a component of the 'stochastic physics scheme'. Multiply tendencies by a number between 0.5 and 1.5, selected randomly with a uniform pdf. 10 degree lat/lon box correlation
- **Lin and Neelin (2000)** allow random CAPE fluctuations in the Betts-Miller convection parametrization. **Lin and Neelin (2002)** use observed rainfall pdf to adapt existing parametrization
- **Plant and Craig (2008)** – stochastic convection parametrization based on a mix of statistical mechanics theory and conventional equilibrium parametrization.
- **Teixeira and Reynolds (2008)** – perturbed wind and temperature tendencies. No temporal/horizontal spatial correlations.

- **Bowler et al (2008) (used in MOGREPS – the Met Office EPS system)**
 - (i) Stochastic convective vorticity – based on anticyclonic lens/meso-vortex model of Gray and Shutts (2002)**
 - (ii) Random parameters – vary entrainment rate and CAPE time scale as an autoregressive process in time**

- **Shutts (2005) Cellular Automaton Backscatter Scheme (CABS) – includes a convective component to return KE generated by buoyancy to larger scales**

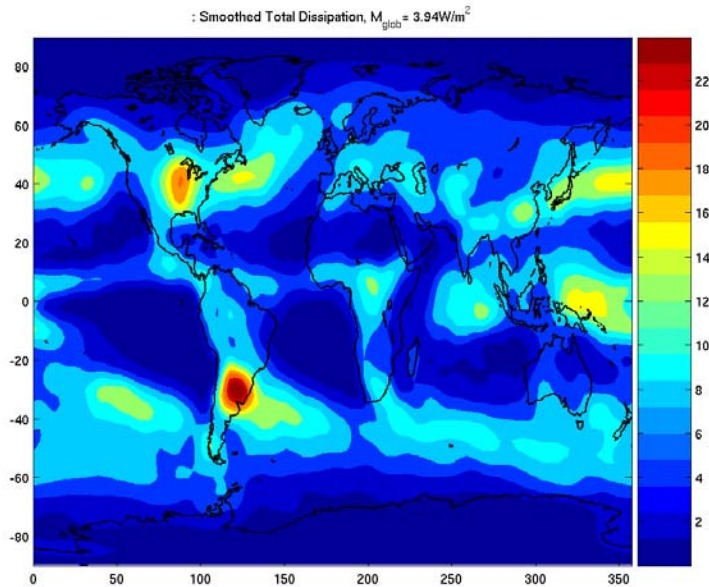
- **Berner et al (2008) - spectral backscatter scheme (adaption of CABS)**

ECMWF Spectral Backscatter Scheme Berner et al, 2008

Rationale: A fraction of the dissipated energy is scattered upscale and acts as streamfunction forcing for the resolved-scale flow (LES, CASBS: Shutts and Palmer 2004, Shutts 2005); **New:** spectral pattern generator

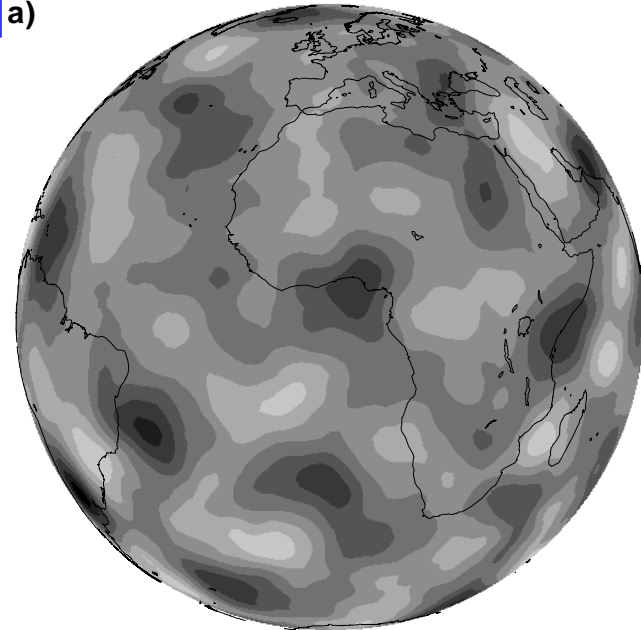
$$\Delta \psi^* \propto \sqrt{D} \psi' \quad \text{a)}$$

D



Total Dissipation rate from numerical dissipation, convection, gravity/mountain wave drag.

ψ'



Spectral Markov chain: temporal and spatial correlations prescribed

Spectral Backscatter scheme

Assume a streamfunction perturbation in **spherical harmonics** representation

$$\psi'(\phi, \lambda) = \sum_{n=0}^N \sum_{m=-n}^n \psi_n^{lm}(t) P_{n,m}(\mu) e^{im\lambda}$$

Assume furthermore that each coefficient evolves according to the **spectral Markov process**

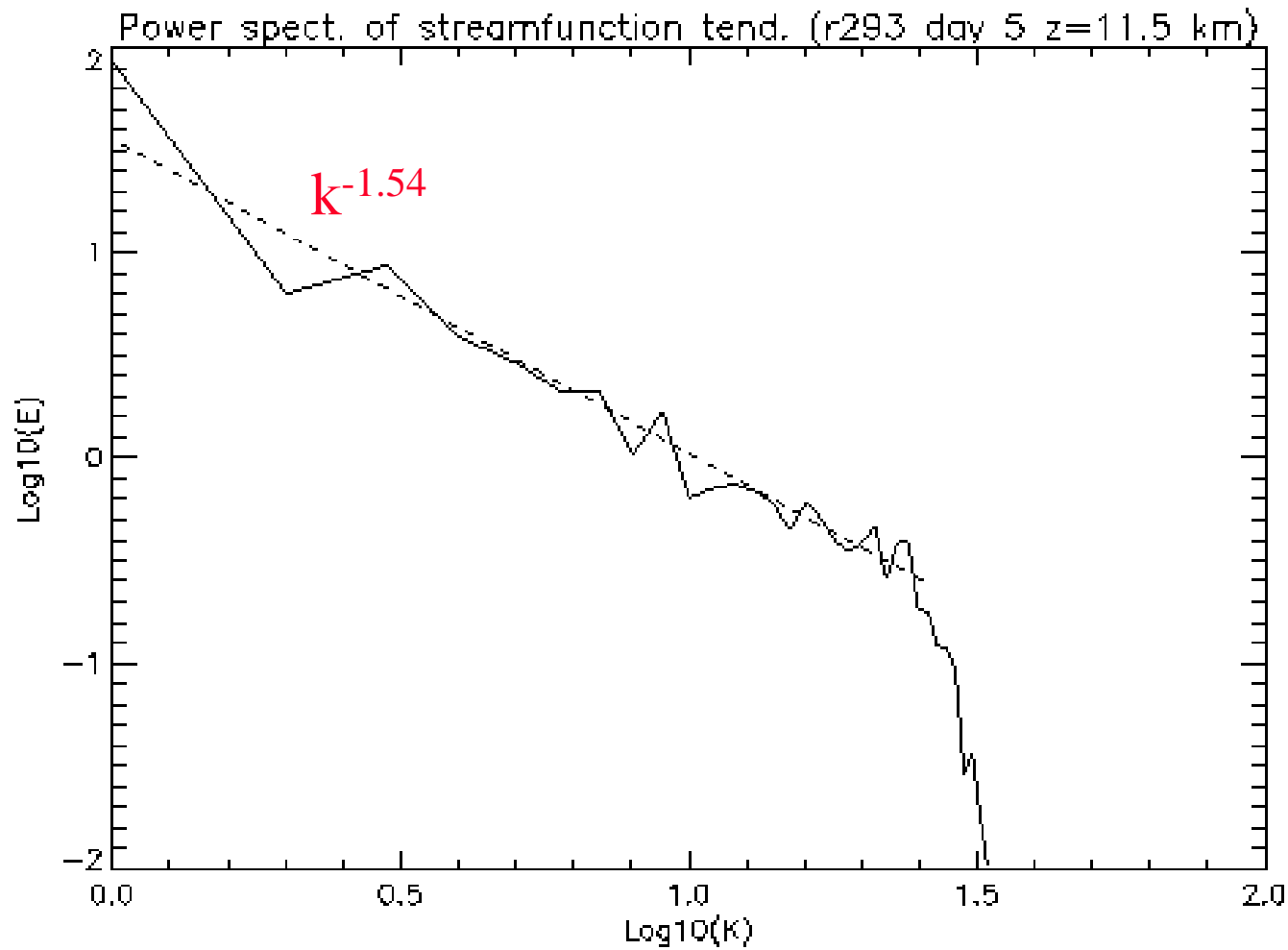
$$\psi_n^{lm}(t+1) = (1 - \alpha)\psi_n^{lm}(t) + g_n \sqrt{\alpha} \epsilon(t)$$

Find the wavenumber dependent noise amplitudes $g_n = b n^p$

so that prescribed kinetic energy dE is injected into the flow

$$b_n = \left(\frac{4\pi a^2 \alpha}{\sigma_z \Gamma} dE' \right)^{\frac{1}{2}} \quad \text{with} \quad \Gamma = \sum_{n=n_1}^{n_2} n(n+1)(2n+1)n^{2p}$$

Power spectrum of coarse-grained streamfunction forcing

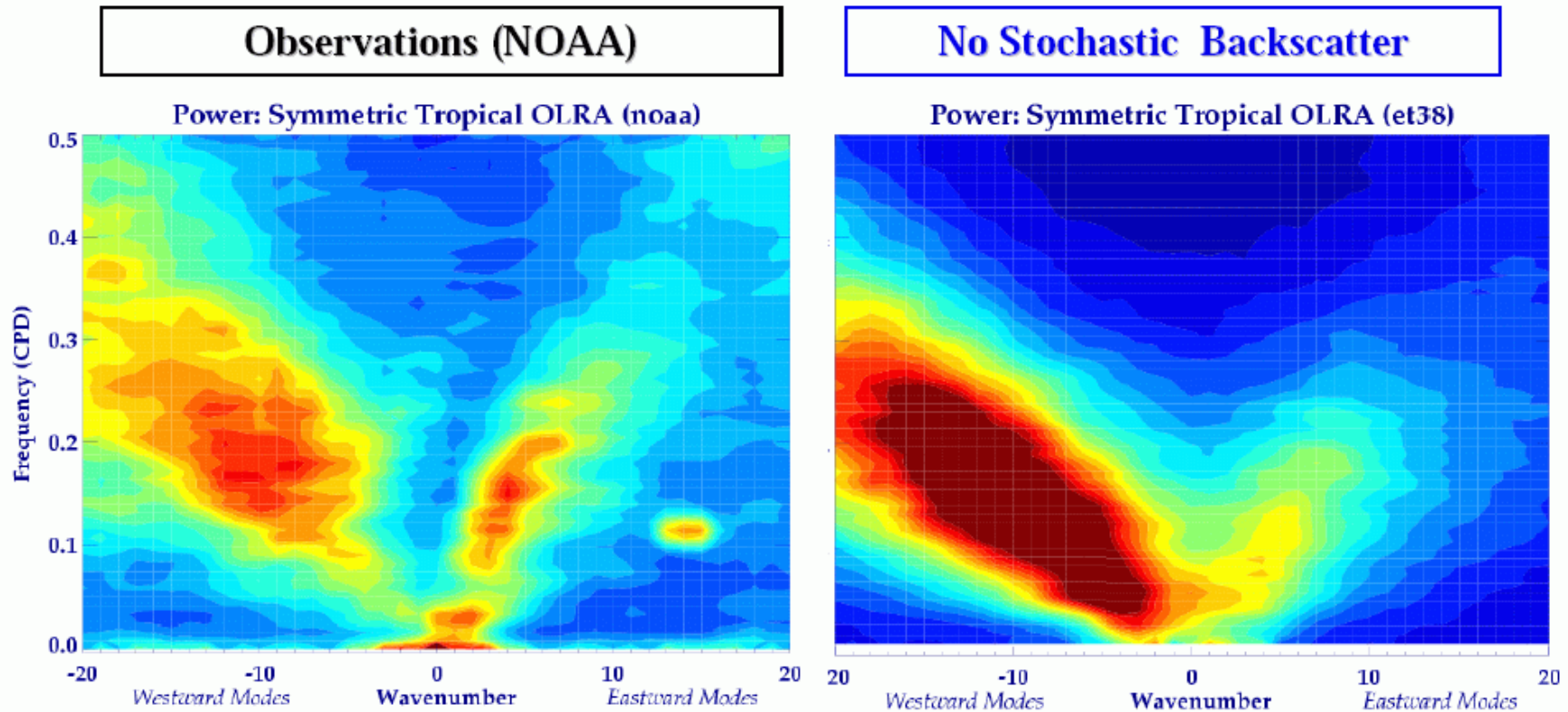


Wavenumber-Frequency Spectrum

Slide from Judith
Berner

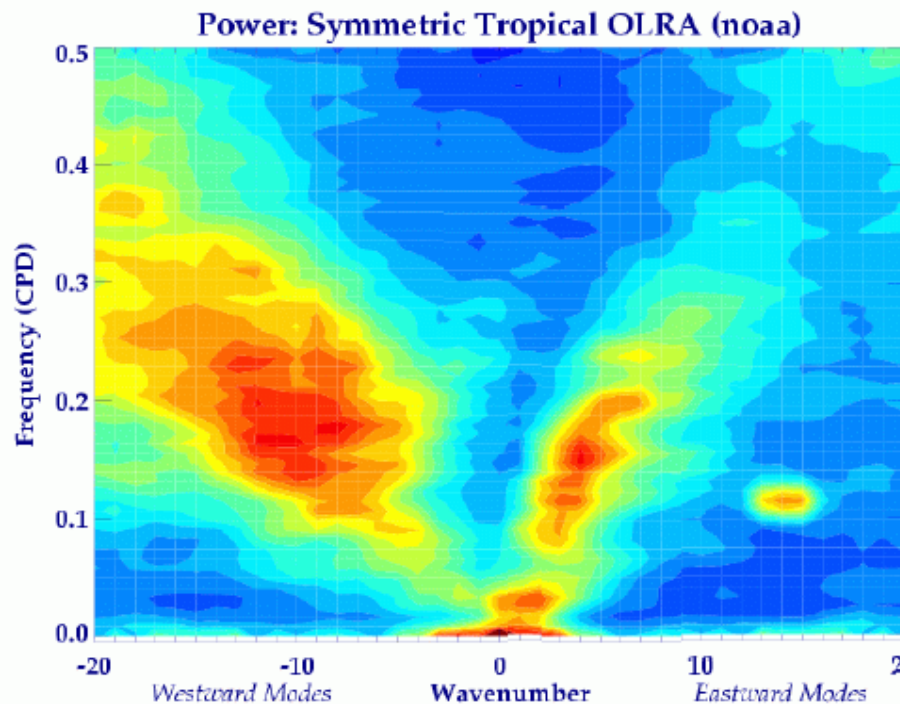
Symmetric part, background removed
(after Wheeler and Kiladis, 1999)

cy31r1

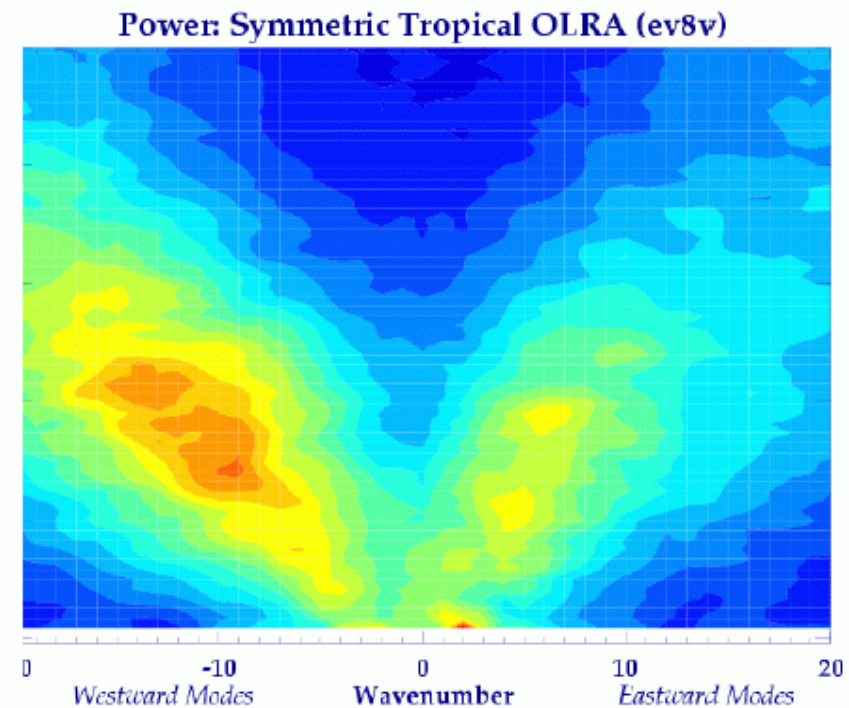


Improvement in Wavenumber-Frequency Spectrum

Observations (NOAA)



Stochastic Backscatter



❖ Backscatter scheme reduces erroneous westward propagating modes

Summary

- **Convection is a multi-scale phenomenon**
- **Convective mass fluxes may generate mesoscale PV anomalies and associated balanced flow structures (e.g. lens and front)**
- **Convective forcing at the near-gridscale is a non-equilibrium phenomenon**
- **Stochastic methods are desirable**
- **Must calibrate these methods using CRMs**
e.g. **CASCADE project**